



ON THE STUDY OF WAVE- PARTICLE INTERACTION EFFECTS IN A DUSTY PLASMA

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Outline

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- Aim
- Physical Model
- Linear Analysis
- Non-linear Analysis
- Discussion
- Conclusion

Introduction

The wave particle interaction plays a crucial role in the structure and dynamics of various space and laboratory plasma system

.... *Why?*

➤ Plasma controlled by long-range electromagnetic waves .

➤ Applications

(space plasmas) formation magneto pause boundary

(astrophysics plasmas) charged particle acceleration in relativistic regimes

(laboratory plasmas) plasma heating in magnetic fusion plasmas

- **What happens between the charged particle and the wave in the plasma?**

1. $(V_T > V_{Ph})$

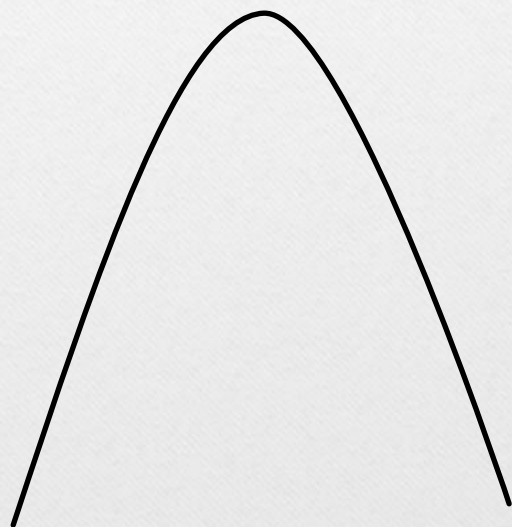
Passing particles (Free particle)

2. $(V_T \leq V_{Ph})$

Resonance particles

- **(Resonance Trapped particle)**
- **(Resonance Reflected particle)**

WPIs



Phase-space of Resonant charges

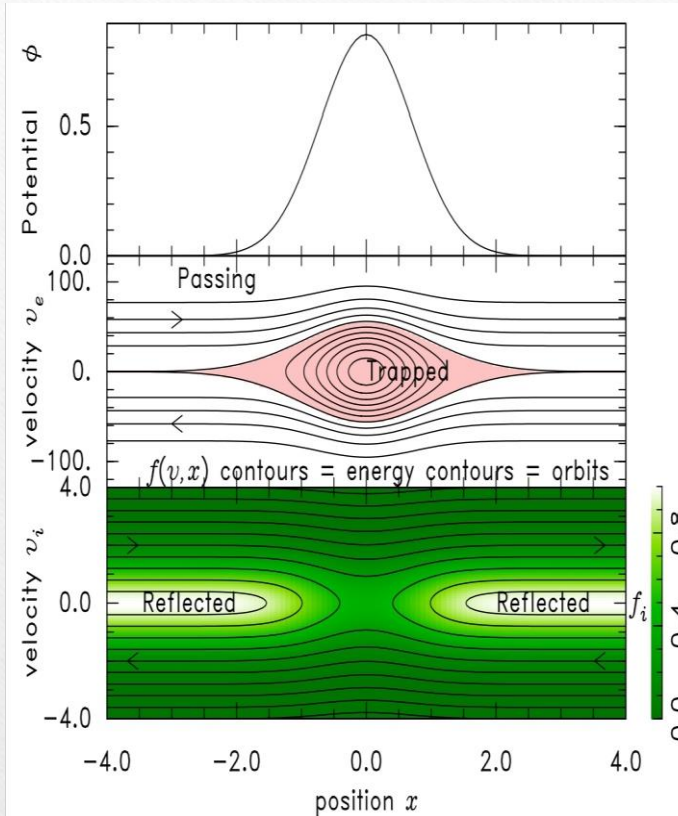


FIG. 1. Schematic of a slow electron hole and the corresponding electron and ion phase-space density contours.

Wave-particle interaction scenario with two different charged particles (SEADL,SIADL)

Distribution Functions

$$f_e^{\kappa_e}(v, \phi) = \frac{\Gamma[\kappa_e]}{\sqrt{2\pi}(\kappa_e - \frac{3}{2})^{1/2}\Gamma[\kappa_e - \frac{1}{2}]} \begin{cases} (1 + \frac{v^2 - \phi}{\kappa_e - \frac{3}{2}})^{-\kappa_e}, & |v| > 2\phi, \\ (1 + \beta_e \frac{v^2 - \phi}{\kappa_e - \frac{3}{2}})^{-\kappa_e}, & |v| \leq 2\phi. \end{cases}$$

$$f_i^{\kappa_i}(u, \phi) = \frac{\Gamma[\kappa_i]}{\sqrt{2\pi}(\kappa_i - \frac{3}{2})^{1/2}\Gamma[\kappa_i - \frac{1}{2}]} \begin{cases} (1 + \frac{u^2 - \theta(\psi - \phi)}{\kappa_i - \frac{3}{2}})^{\kappa_i}, & |u| > 2\theta(\psi - \phi), \\ (1 + \alpha_i \frac{u^2 - \theta(\psi - \phi)}{\kappa_i - \frac{3}{2}})^{-\kappa_i}, & |u| \leq 2\theta(\psi - \phi). \end{cases}$$

$$f_e^{\kappa_e}(v, \phi) = f_{fe}^{\kappa_e}(v, \phi) + f_{te}^{\kappa_e}(v, \phi), \quad f_i^{\kappa_i}(v, \phi) = f_{fi}^{\kappa_i}(u, \phi) + f_{ri}^{\kappa_i}(u, \phi)$$

Aim

Study of the WPIs (Trapping + Reflection) and Superthermality of two different charges (electron, ion) effects on DAW:

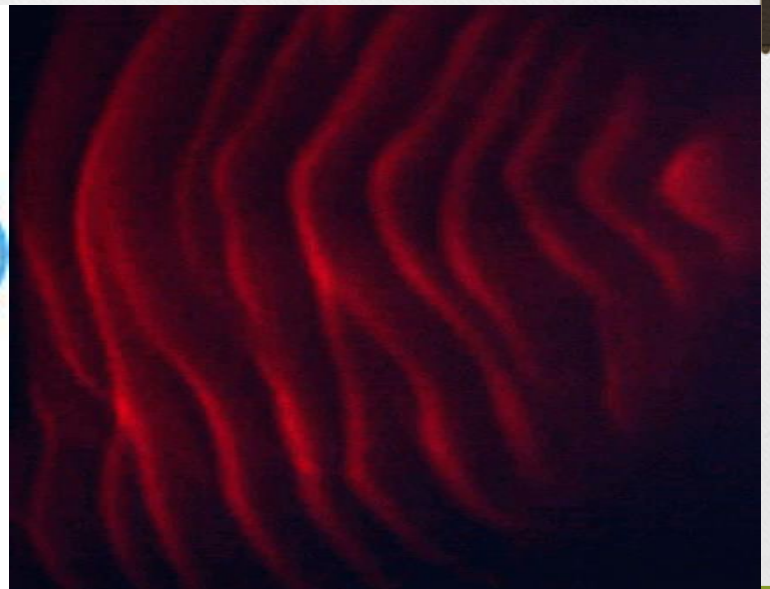
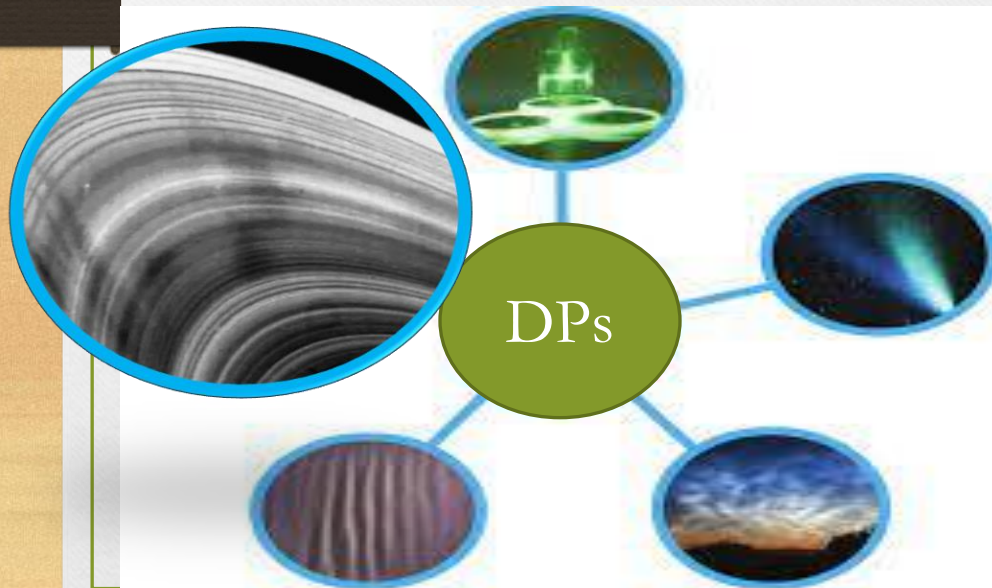
1. Existence
2. No of modes
3. Polarity
4. Profile
5. Propagation speed

WPIs Cases: $\left\{ \begin{array}{l} \text{Case A (TE+RI)} \\ \text{Case B (TI+RE)} \end{array} \right.$

Modelling

Plasma system:

Un-magnetized collision-less adiabatic OPDP system consisting of positive dust, negative dust, nonthermal electrons and nonthermal ions.



Model Equations

$$\frac{\partial n_n}{\partial t} + \frac{\partial(n_n u_n)}{\partial x} = 0$$

$$\frac{\partial u_n}{\partial t} + u_n \frac{\partial u_n}{\partial x} = \frac{\partial \phi}{\partial x} - \sigma_n n_n \frac{\partial n_n}{\partial x},$$

$$\frac{\partial n_p}{\partial t} + \frac{\partial(n_p u_p)}{\partial x} = 0$$

$$\frac{\partial u_p}{\partial t} + u_p \frac{\partial u_p}{\partial x} = -v \frac{\partial \phi}{\partial x} - \sigma_p n_p \frac{\partial n_p}{\partial x},$$

$$\frac{\partial^2 \phi}{\partial x^2} = n_n - \delta_p n_p + \delta_e n_e - \delta_i n_i$$

general mathematical form for the number density

$$n_e \simeq 1 + c_{e1}\phi + c_{e2}\phi^{3/2} + c_{e3}(\psi^{3/2} - (\psi + \phi)^{3/2}) + c_{e4}\phi^2$$

$$n_i \simeq 1 + c_{i1}\phi + c_{i2}(-\phi)^{3/2} + c_{i3}(\psi^{3/2} - (\psi - \phi)^{3/2}) + c_{i4}\phi^2$$

$$c_{e1} = \frac{2\kappa_e - 1}{2\kappa_e - 3},$$

$$c_{e2} = \frac{8\sqrt{2/\pi}(\beta_e - 1)\kappa_e\Gamma[\kappa_e]}{3(2\kappa_e - 3)^{3/2}\Gamma[\kappa_e - \frac{1}{2}]},$$

$$c_{e3} = -\frac{8\sqrt{2/\pi}(\alpha_e - 1)\kappa_e\Gamma[\kappa_e]}{3(2\kappa_e - 3)^{3/2}\Gamma[\kappa_e - \frac{1}{2}]},$$

$$c_{e4} = \frac{4\kappa_e^4 - 1}{2(2\kappa_e - 3)^2}$$

$$c_{i1} = \theta \left(\frac{2\kappa_i - 1}{2\kappa_i - 3} \right),$$

$$c_{i2} = \theta^{3/2} \left(\frac{8\sqrt{\frac{2}{\pi}}(\beta_i - 1)\kappa_i\Gamma[\kappa_i]}{3(2\kappa_i - 3)^{\frac{3}{2}}\Gamma[\kappa_i - \frac{1}{2}]} \right),$$

$$c_{i3} = \theta^{3/2} \left(\frac{8\sqrt{2/\pi}(\alpha_i - 1)\kappa_i\Gamma[\kappa_i]}{3(2\kappa_i - 3)^{3/2}\Gamma[\kappa_i - \frac{1}{2}]} \right),$$

$$c_{i4} = \theta^2 \left(\frac{4\kappa_i^4 - 1}{2(2\kappa_i - 3)^2} \right)$$

Linear Wave Analysis

$$LDR \quad \omega^4 + a_1 \omega^2 + a_0 = 0$$

Where;

$$a_1 = - \left(\sigma_n + \sigma_p + \frac{1 + \nu \delta \rho}{\kappa^2 + c_e \delta e - c_i \delta i} \right) \kappa^2,$$

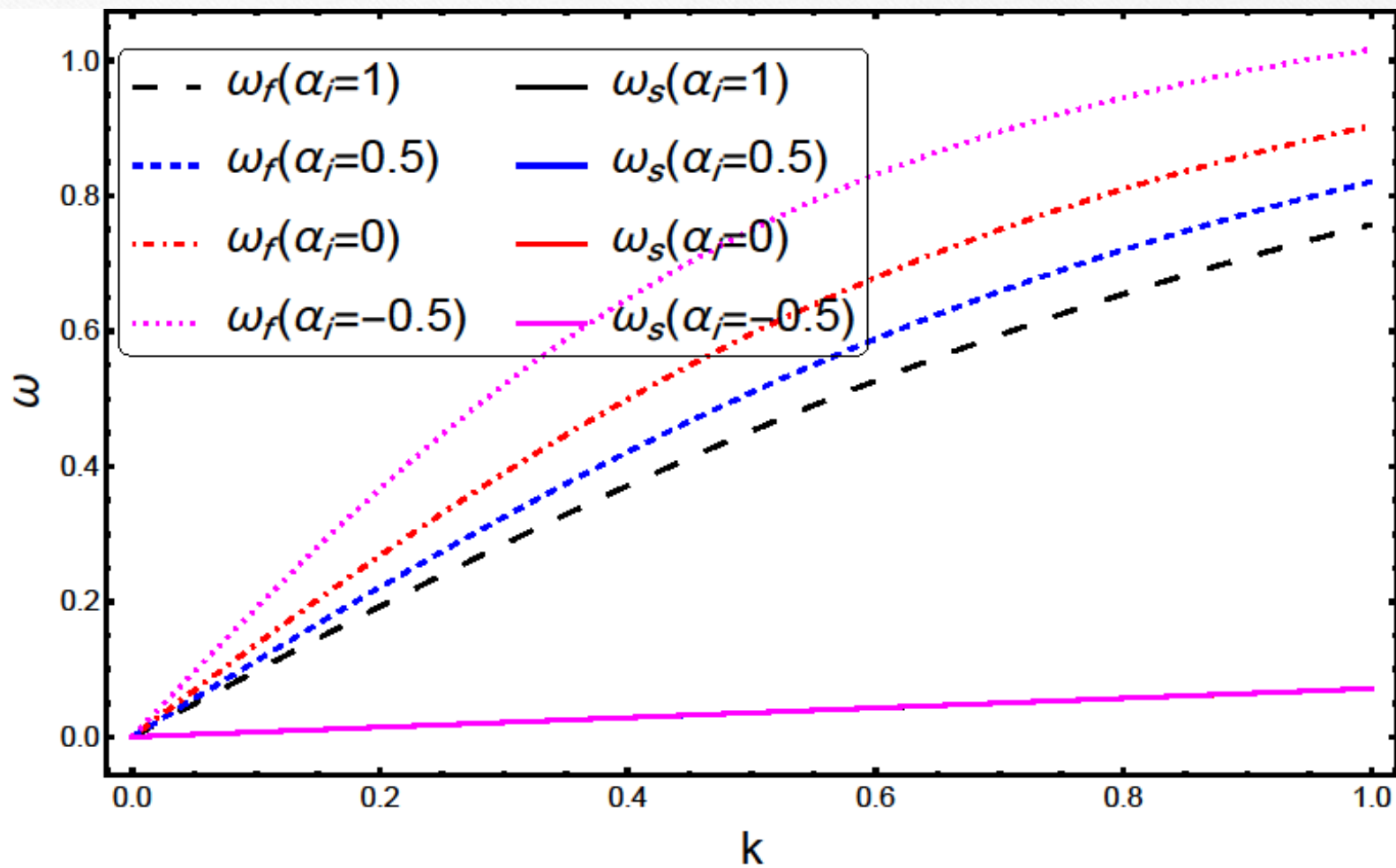
$$a_0 = \left(\sigma_n \sigma_p + \frac{\sigma_p + \nu \delta \rho \sigma_n}{\kappa^2 + c_e \delta e - c_i \delta i} \right) \kappa^4$$

$$c_e = \left(c_{e1} - \frac{3}{2} c_{e3} \psi^{1/2} \right),$$

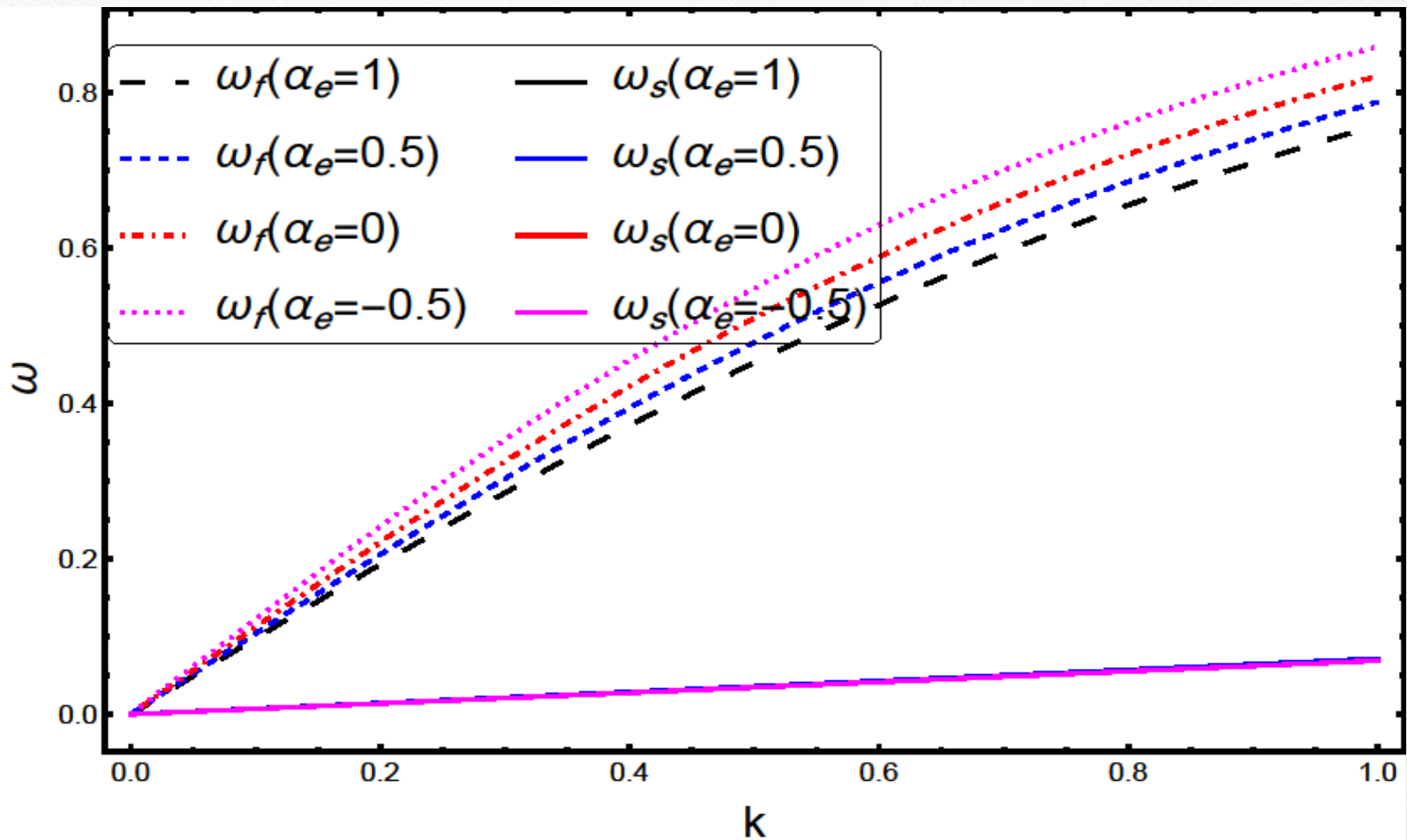
$$c_i = \left(c_{i1} + \frac{3}{2} c_{i3} \psi^{1/2} \right).$$

Case A

($\alpha_e = 1$)



Case B ($\alpha_i = 1$)



Non-Linear Wave Analysis (RPT)

Stretched coordinates $\xi = \epsilon^{1/2}(\mathbf{x} - \lambda t), \quad \tau = \epsilon^{3/2}t.$

Expanding the independent in powers of ϵ

$$n = 1 + \epsilon n^{(1)} + \epsilon^{3/2} n^{(2)} + \epsilon^2 n^{(3)} + \dots,$$

$$u = \epsilon u^{(1)} + \epsilon^{3/2} u^{(2)} + \epsilon^2 u^{(3)} + \dots,$$

$$\phi = \epsilon \phi^{(1)} + \epsilon^{3/2} \phi^{(2)} + \epsilon^2 \phi^{(3)} + \dots$$

To the next higher order of ϵ , we obtain a group of Eqs., which lead to the following modified Schamel-Kdv-equation(SKdV)

$$\frac{\partial \phi}{\partial \tau} + \left(A_1 \phi + A_2 \phi^{1/2} + A_3 (-\phi)^{1/2} \right) \frac{\partial \phi}{\partial x} + B \frac{\partial^3 \phi}{\partial x^3} = 0$$

$$A_1 = \left(-\frac{3\lambda^2 + \sigma_n}{(\lambda^2 - \sigma_n)^3} + \frac{3\lambda^2 + \sigma_p}{(\lambda^2 - \sigma_p)^3} \nu^2 \delta_p - 2\delta_e \left(c_{e4} - \frac{3}{8} c_{e3} \psi^{-1/2} \right) + 2\delta_i \left(c_{i4} - \frac{3}{8} c_{i3} \psi^{-1/2} \right) \right) B$$

$$A_2 = -\frac{3}{2} \delta_e c_{e2} B,$$

$$A_3 = -\frac{3}{2} \delta_i c_{i2} B,$$

$$B = \frac{1}{2\lambda} \left(\frac{1}{(\lambda^2 - \sigma_n)^2} + \frac{\nu \delta_p}{(\lambda^2 - \sigma_p)^2} \right)^{-1}$$

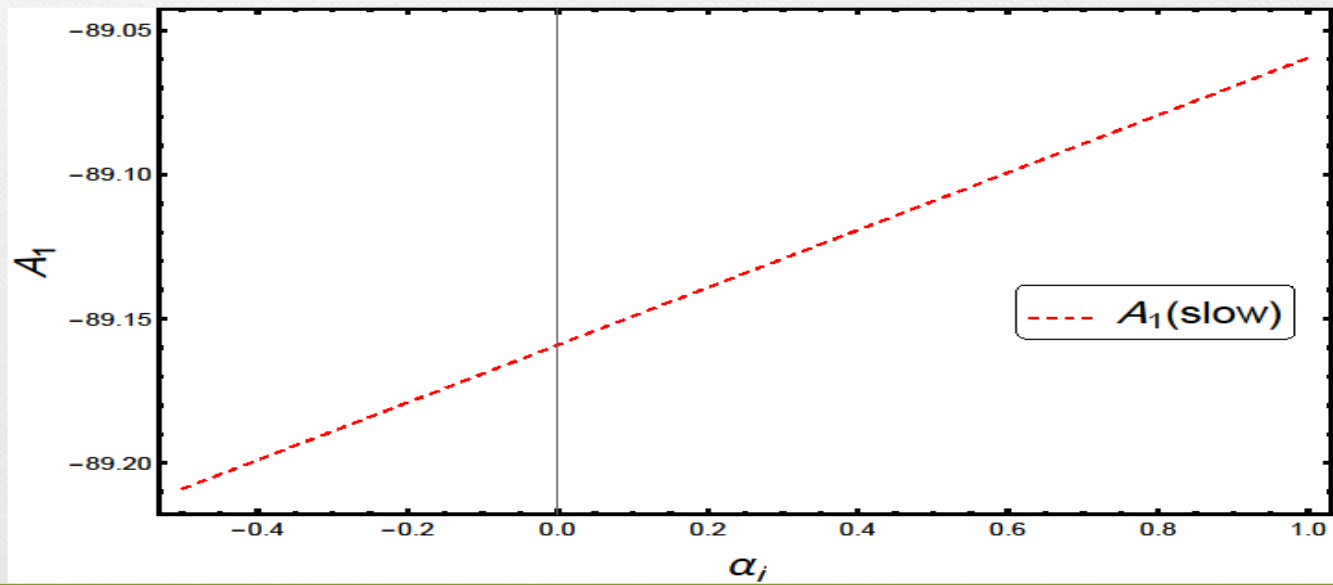
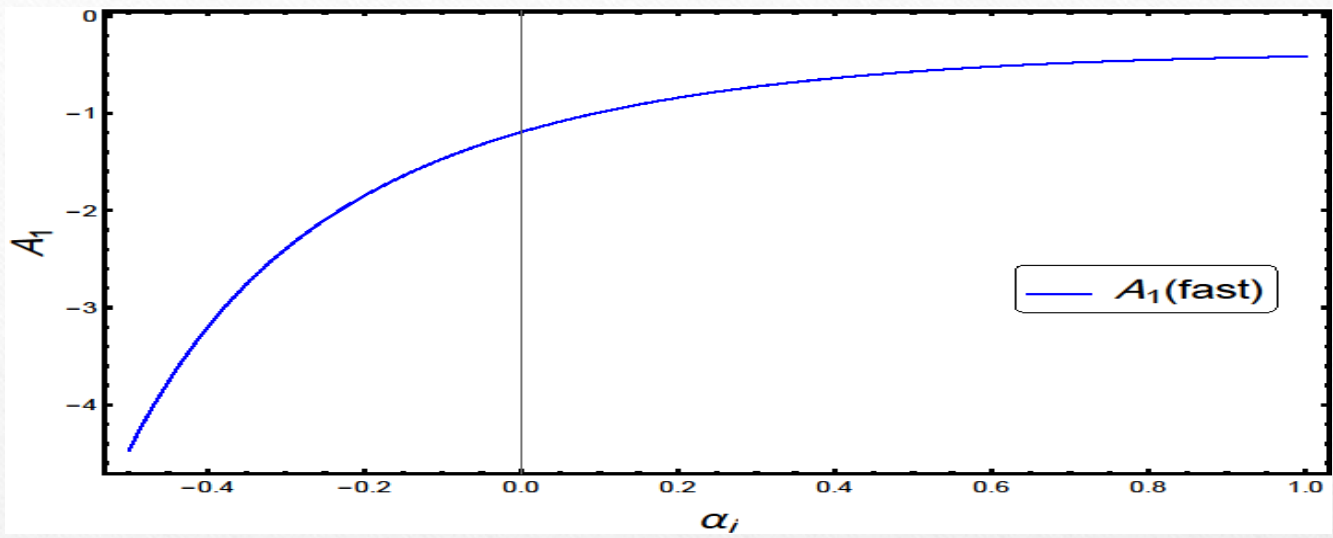
Case A

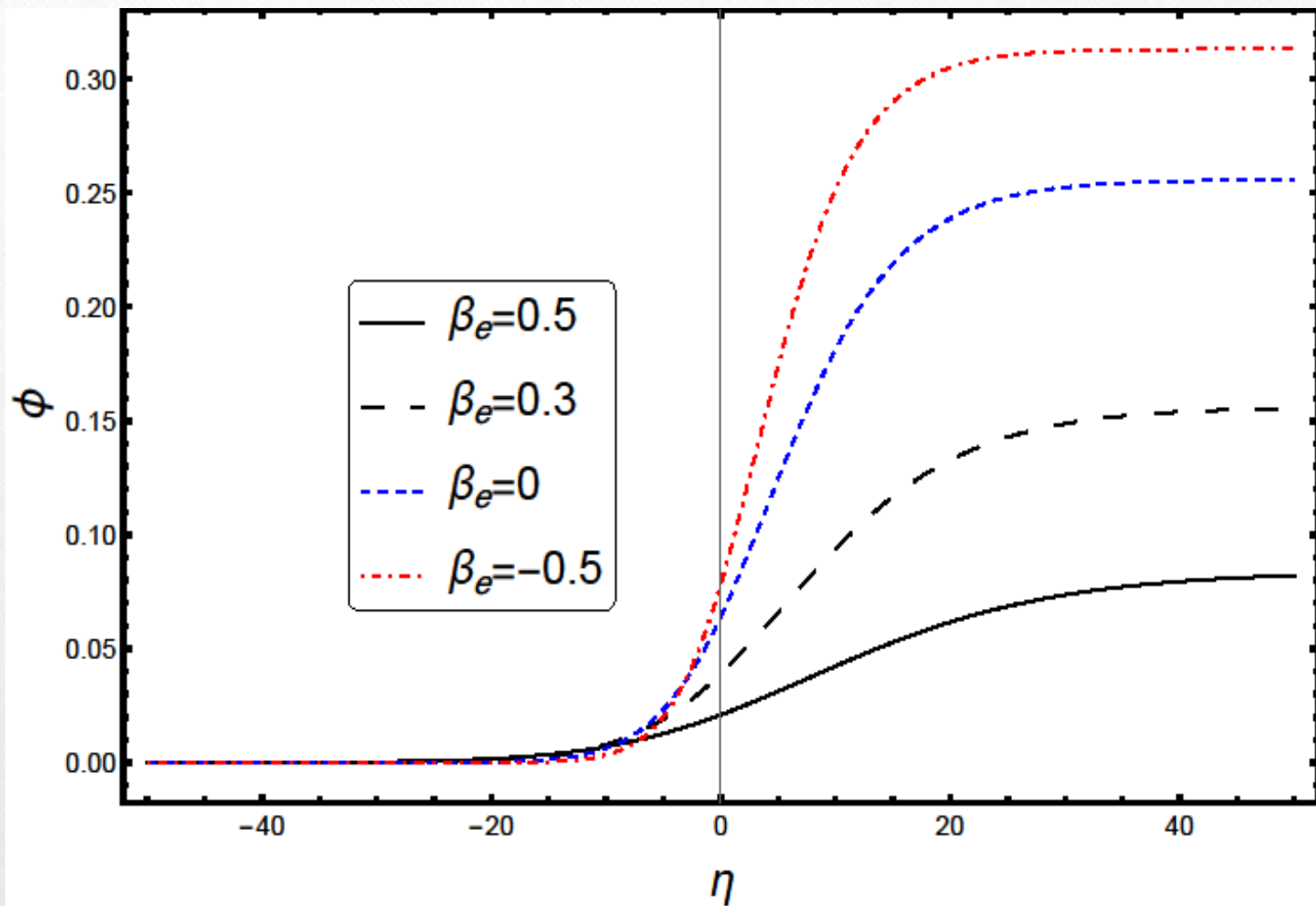
$$(\beta_e < 1, \alpha_i < 1; \beta_i = 1, \alpha_e = 1)$$

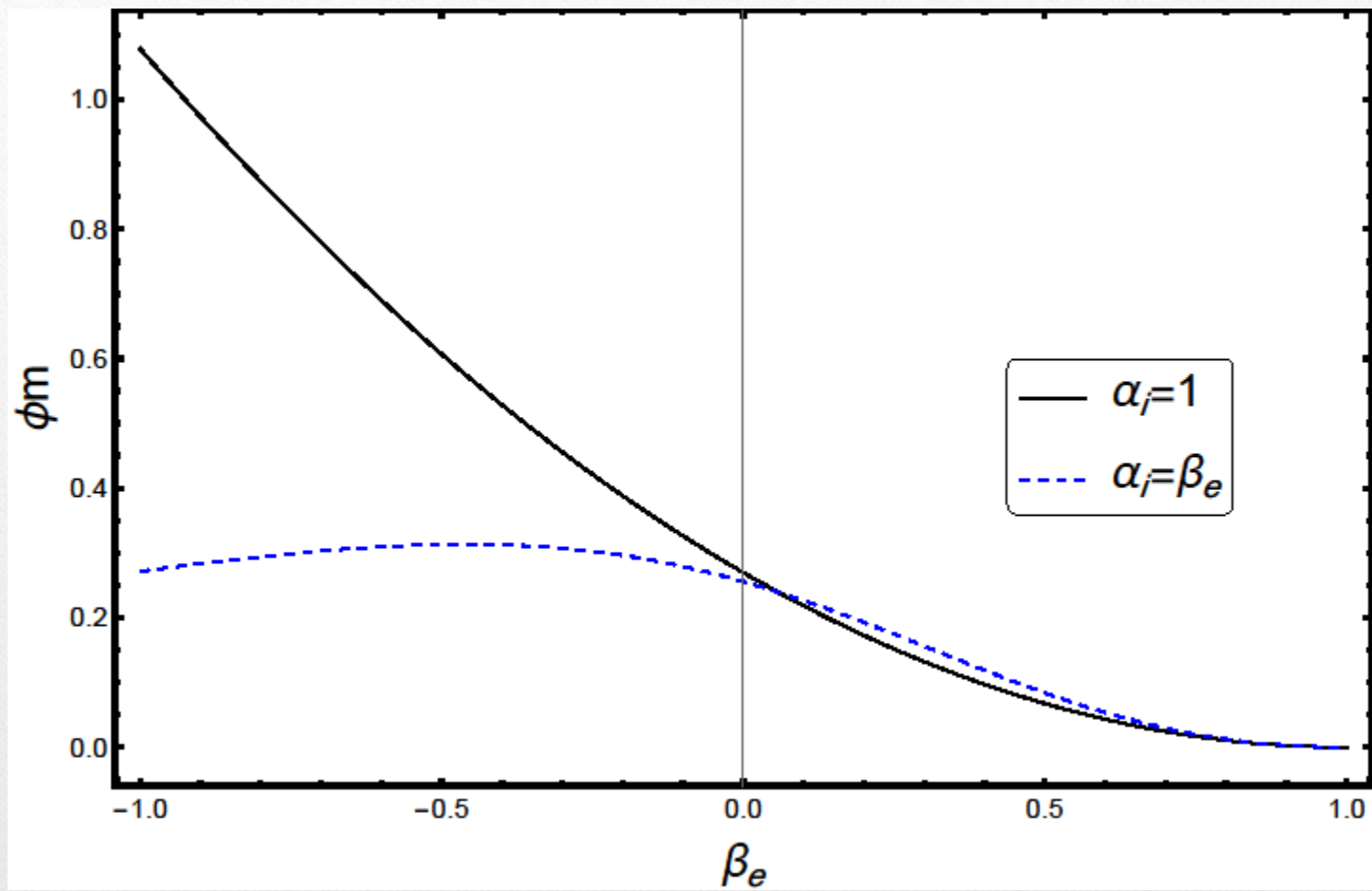
$$\frac{\partial \varphi}{\partial \tau} + \left(A_1 \varphi + A_2 \varphi^{\frac{1}{2}} \right) \frac{\partial \varphi}{\partial \xi} + B \frac{\partial^3 \varphi}{\partial \xi^3} = 0$$

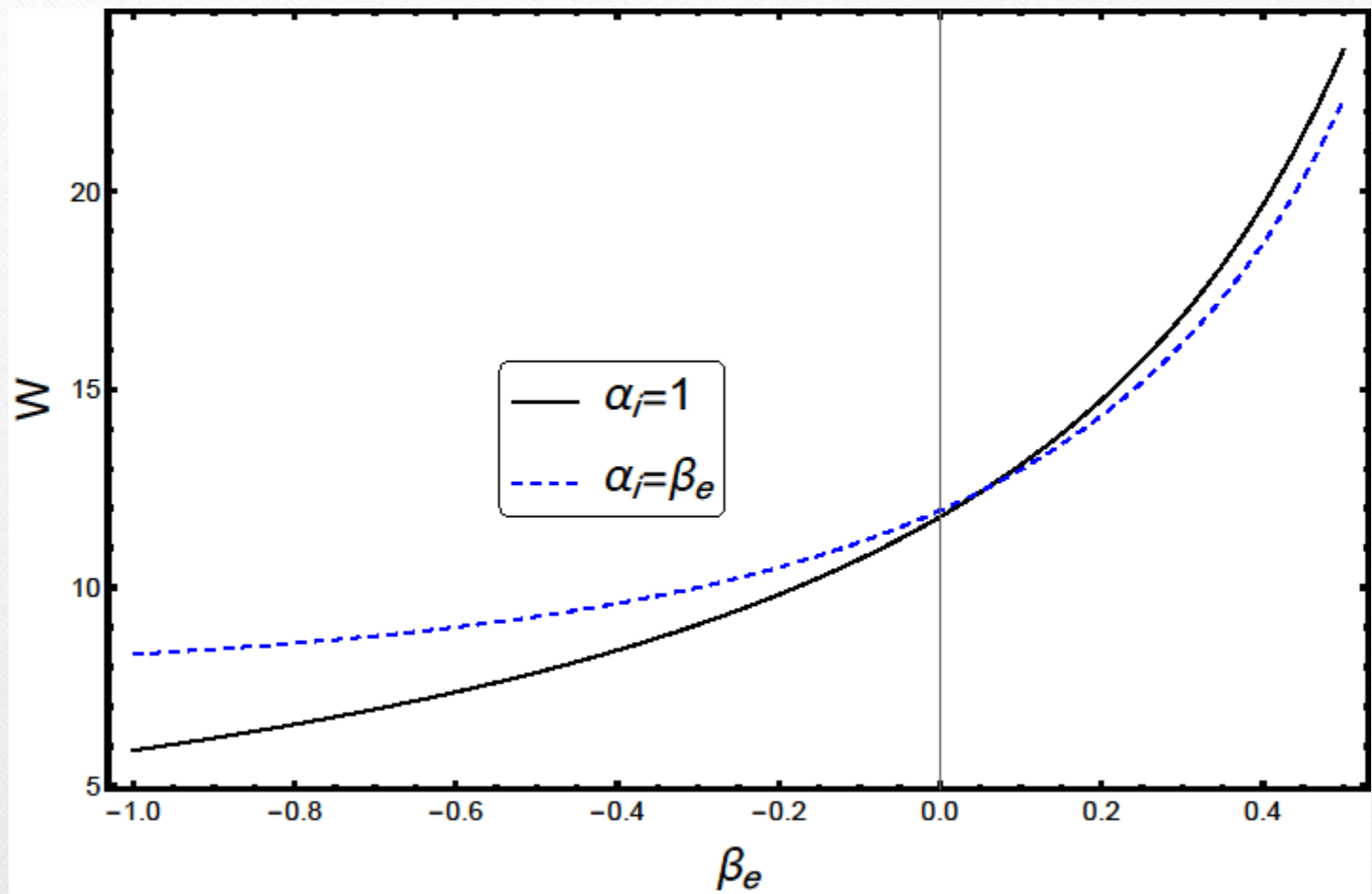
$$\varphi = \varphi_m \left(1 \pm \tanh \left[\frac{\xi - M\tau}{W} \right] \right)^2$$

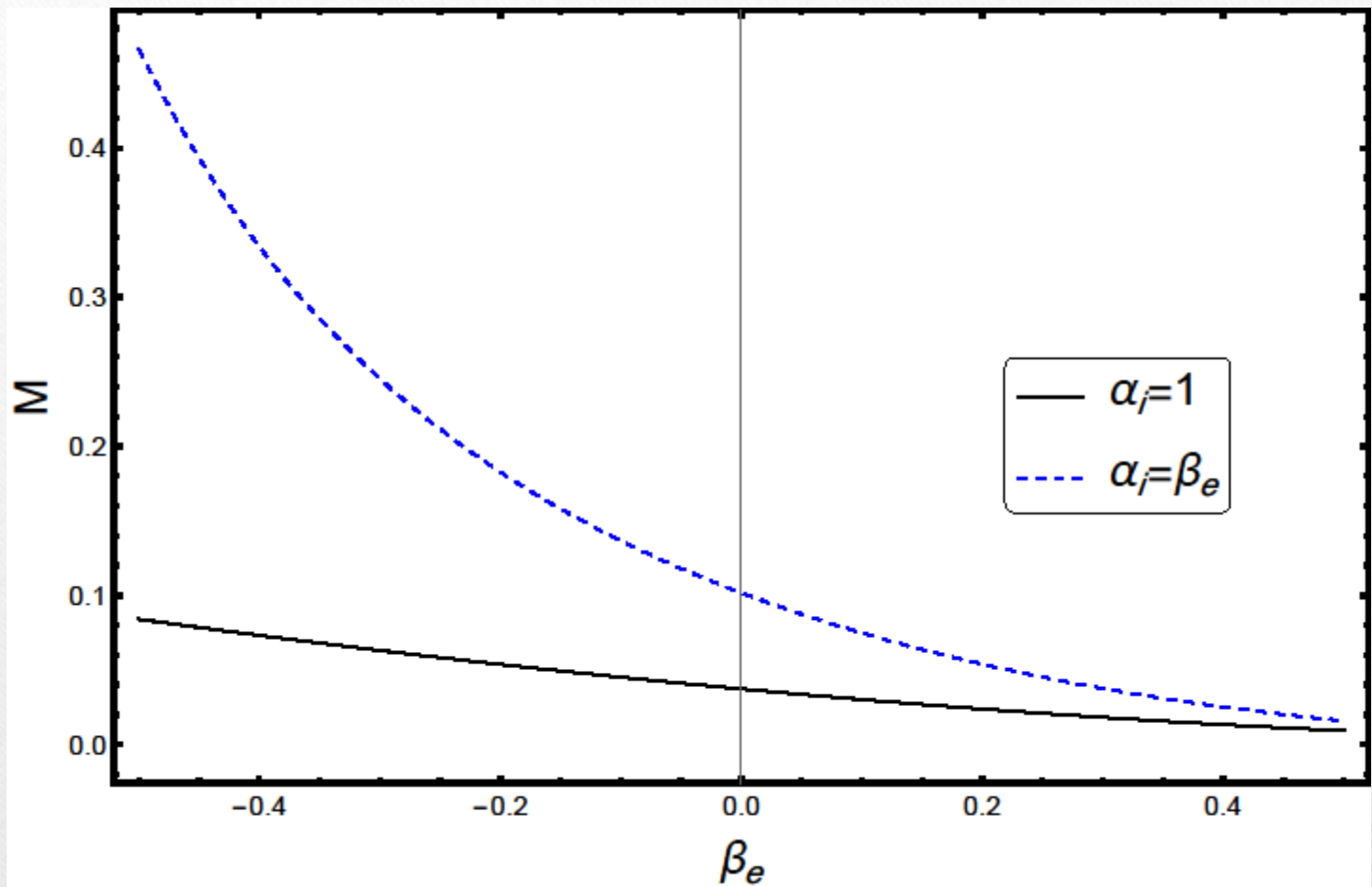
$$\varphi_m = \frac{4A_2^2}{25A_1^2}, \quad W = \frac{5}{A_2} \sqrt{-3BA_1}$$











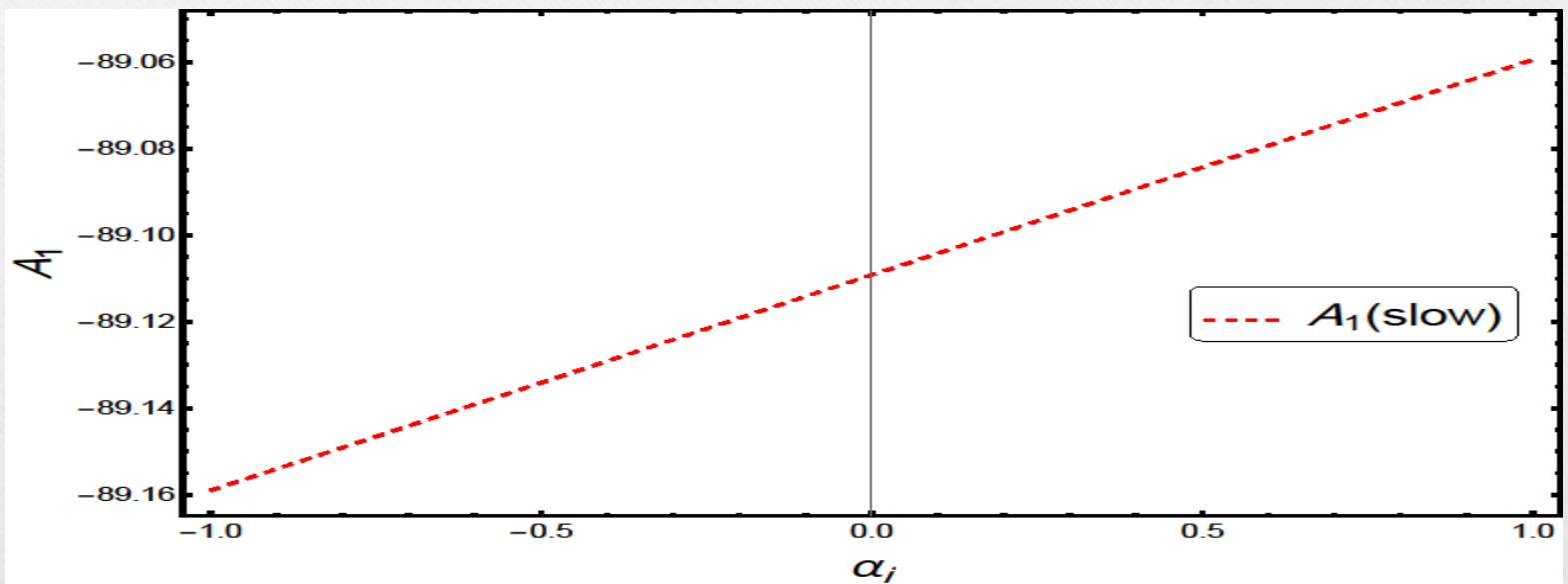
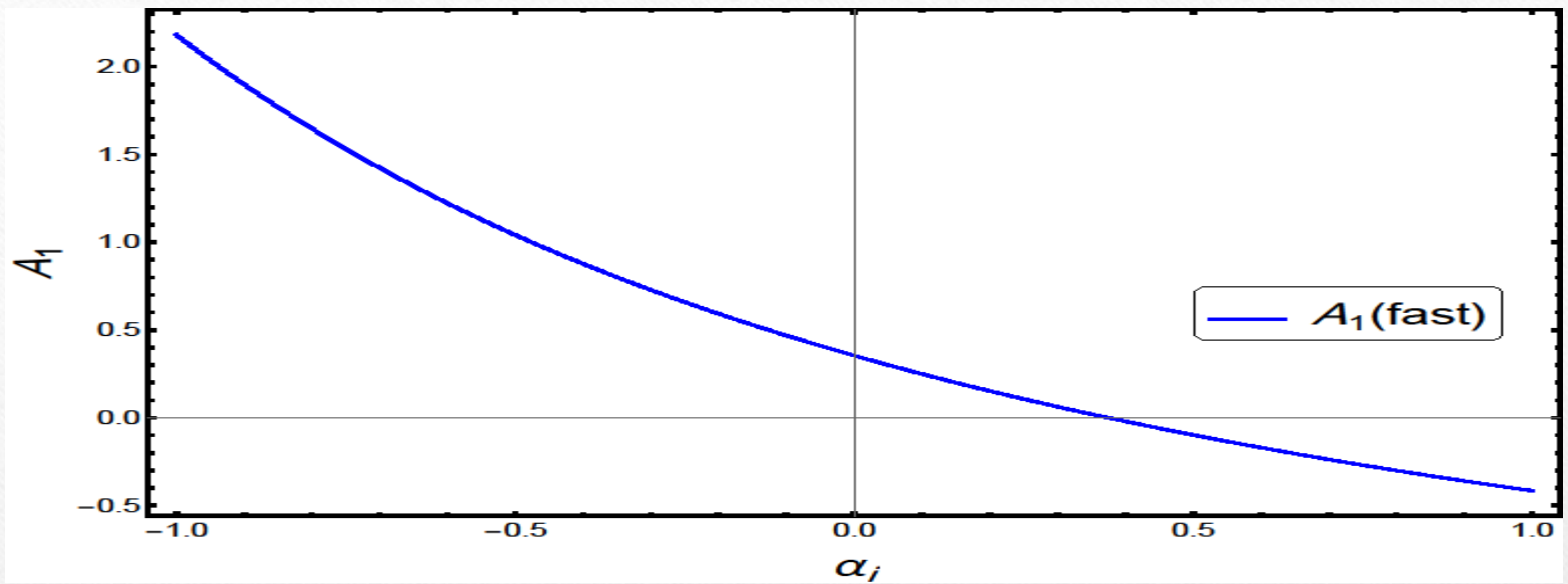
Case B

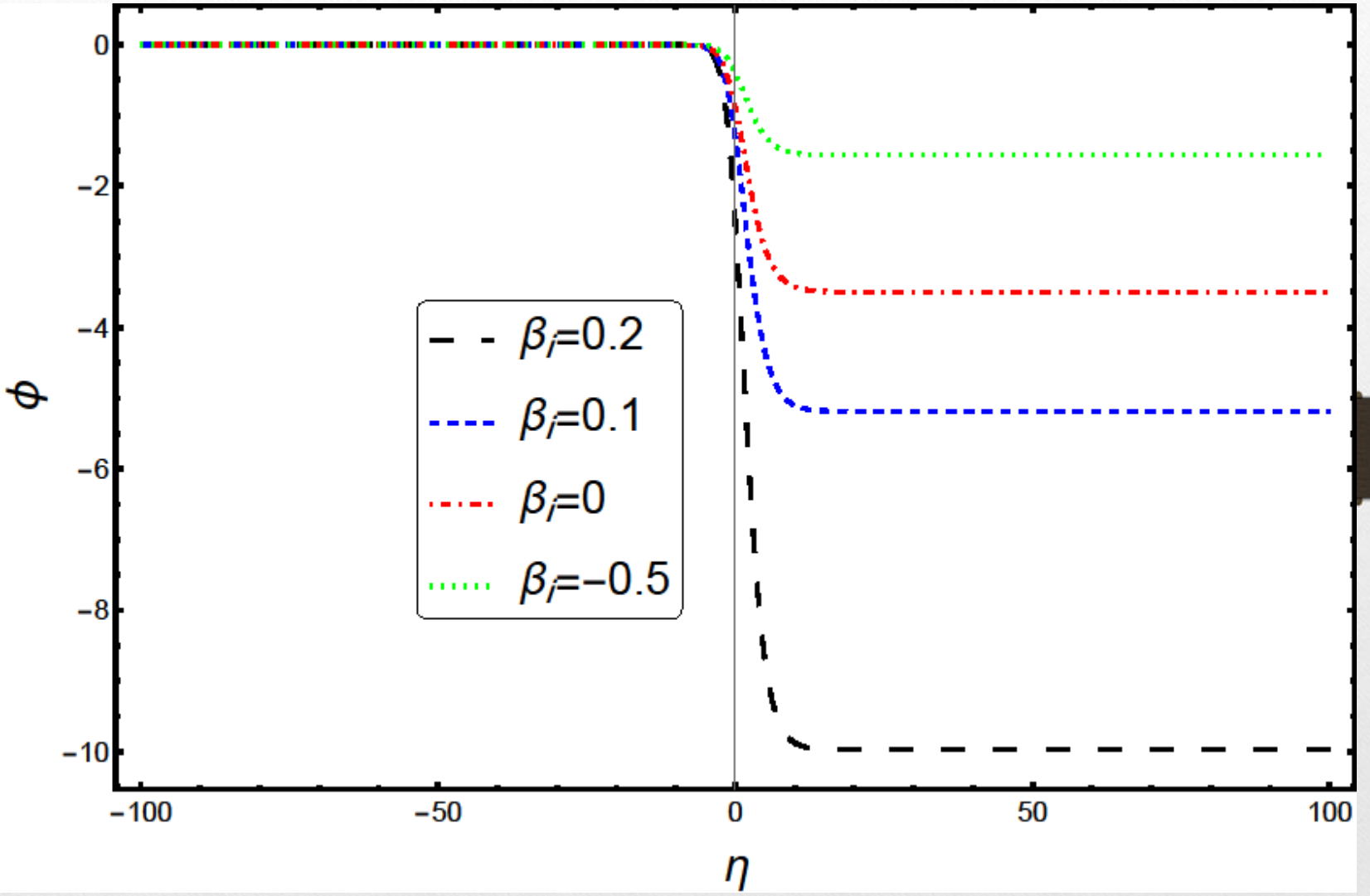
$(\beta_i < 1, \alpha_e < 1; \beta_e = 1, \alpha_i = 1)$

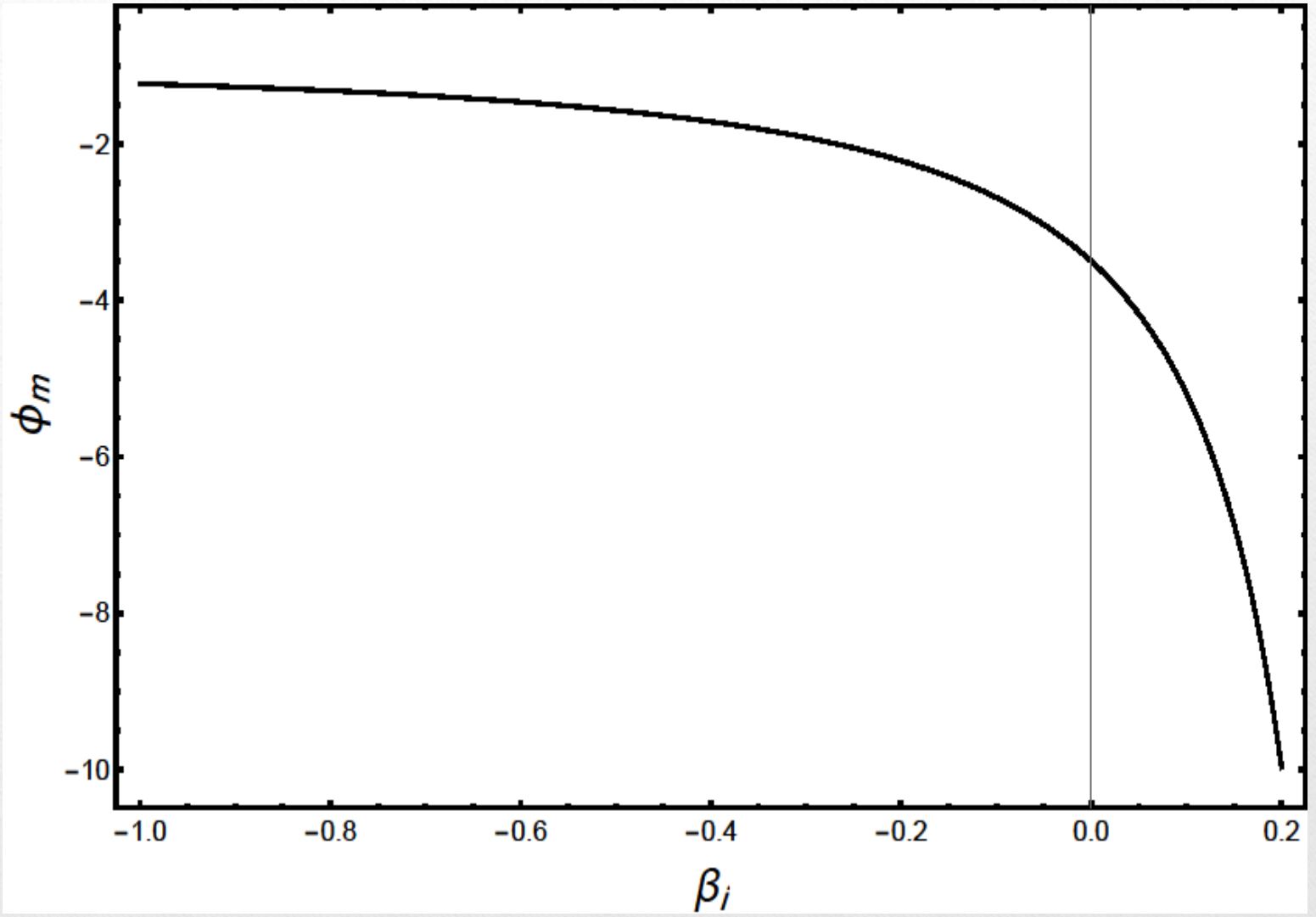
$$\frac{\partial \varphi}{\partial \tau} + \left(A_1 \varphi + A_3 (-\varphi)^{\frac{1}{2}} \right) \frac{\partial \varphi}{\partial \xi} + B \frac{\partial^3 \varphi}{\partial \xi^3} = 0$$

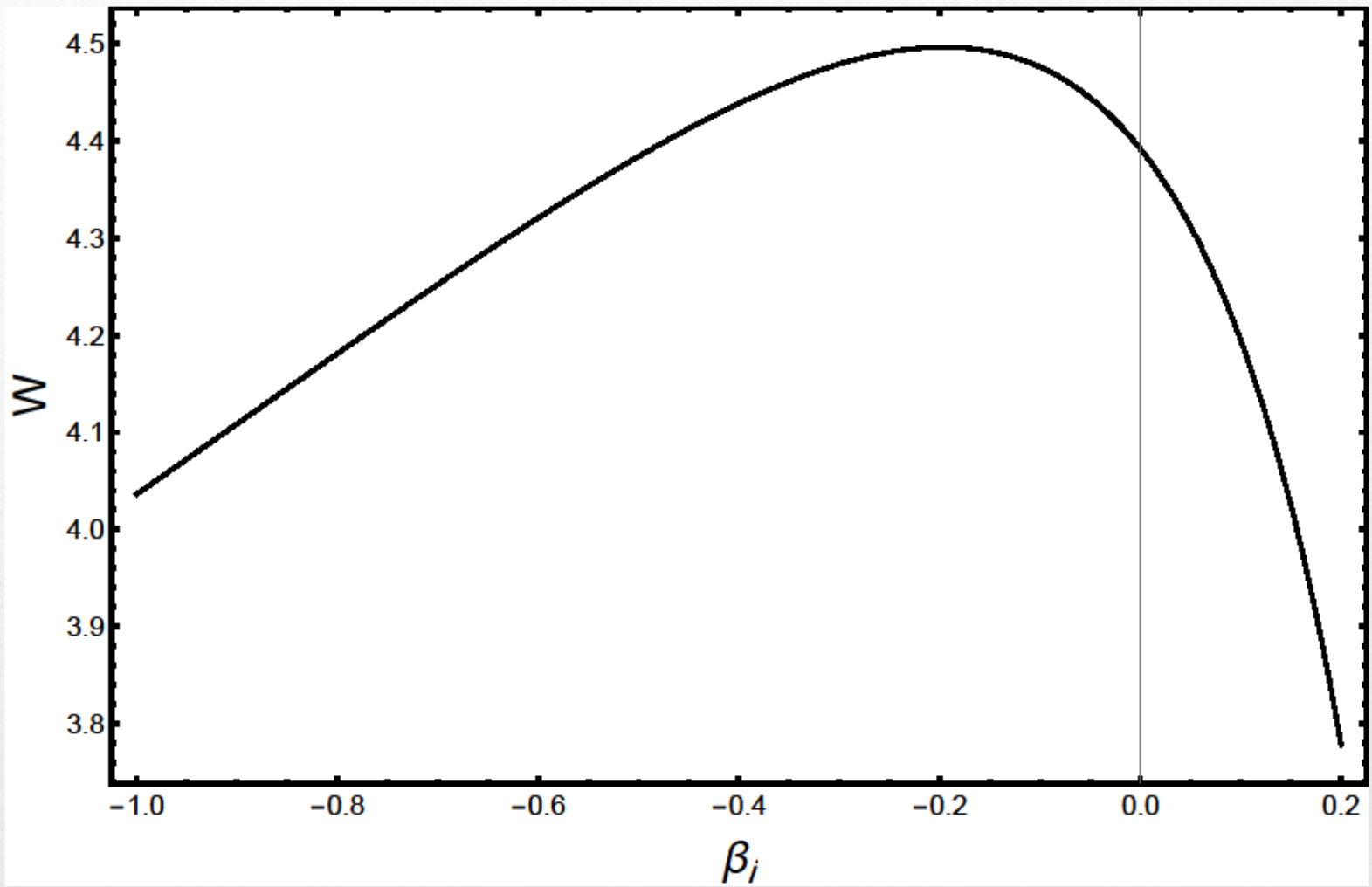
$$\varphi = \varphi_m \left(1 \pm \tanh \left[\frac{\xi - M\tau}{W} \right] \right)^2$$

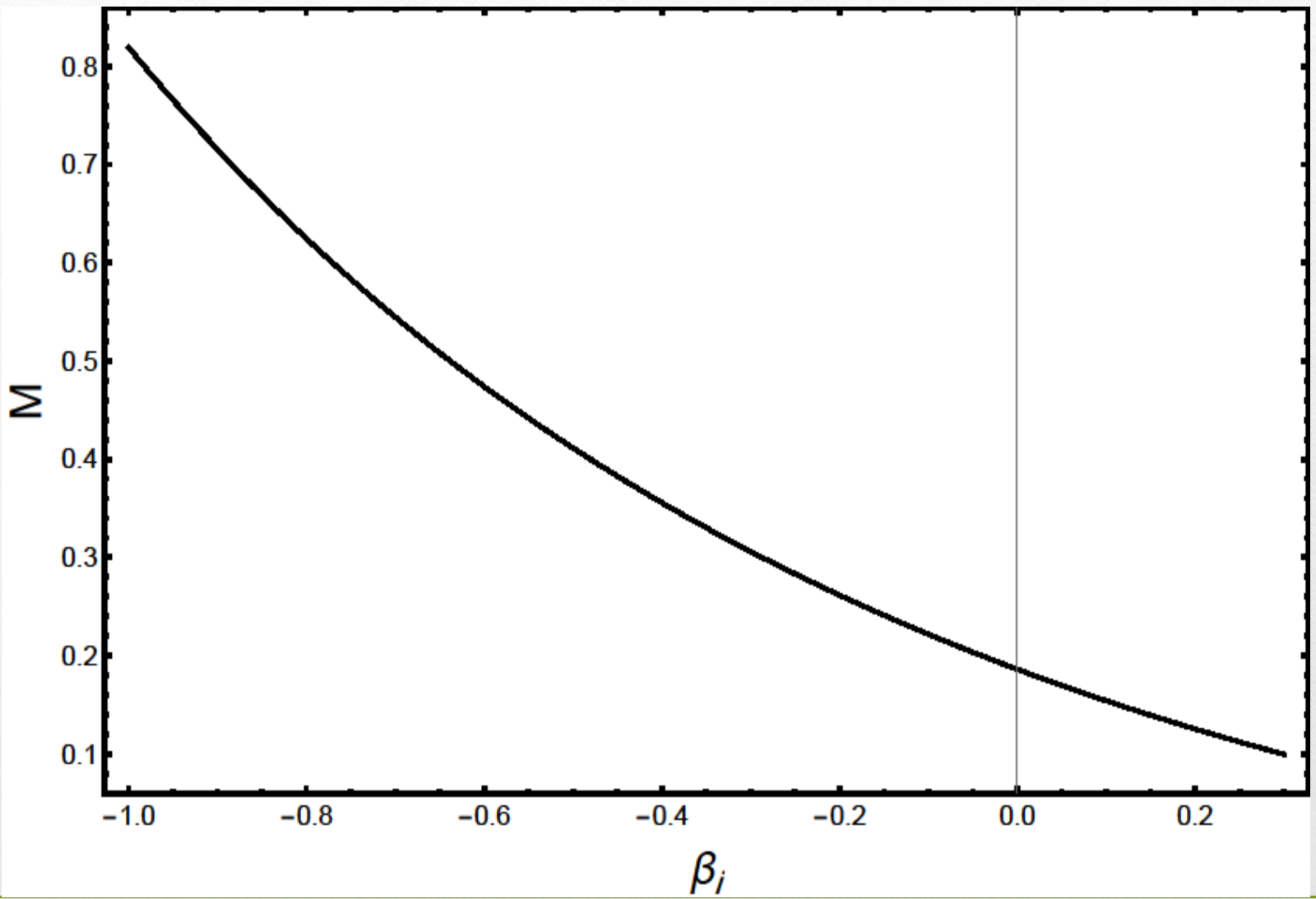
$$\varphi_m = -\frac{4A_3^2}{25A_1^2}, \quad W = \frac{5}{A_3} \sqrt{3BA_1}$$











Conclusions

Case A

Case B

1. Existence	DADL	DADL
2. No. modes	fast, slow	fast
3. Polarity	Positive	Negative
4. Profile	Smaller Amp., Wider width	Larger Amp., Wider width
5. Speed	Slower Speed	Slower Speed

Thanks