

Introduction to Plasma Physics

* Content:

① Historical background

② plasma definition

③ plasma parameters

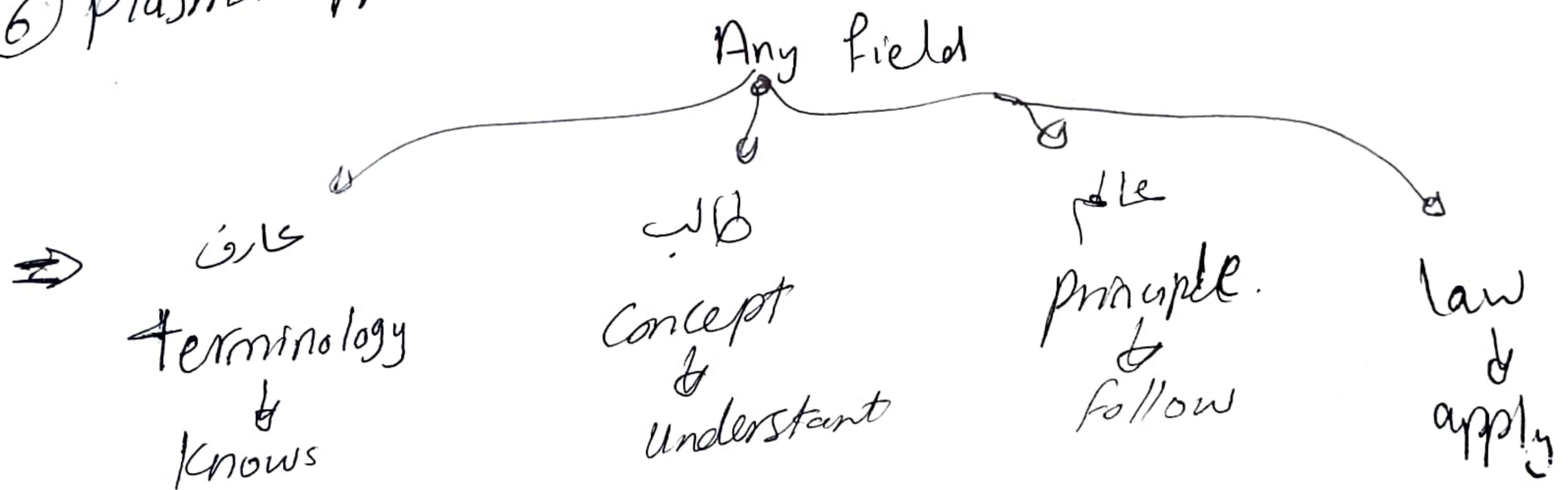
④ plasma criteria

⑤ plasma classification

⑥ plasma applications

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* Plasma history:

- when the blood is cleared of its various corpuscles there remains a transparent liquid which was named plasma (In Greek: jelly or moldable substance)
- 1917: Langmuir used it to describe the ionized gas.
 blood plasma carries white and red corpuscles
 radio transmiss. Ionosphere
current rectification: Gasous electron tube.
 positive charges.
negative charges.
- 1940: Alfven and MHD for Sakur Flowers and Wind.
 MHD magnetic Reconnection.
dynamo theory.
- 1952: Controlled thermonuclear (hydrogen bomb).
- 1958: Van Allen radiation belts (space plasma).
- 1970: laser plasma physics $\rightarrow$ plasma processing

Plasma definition: (Qualitatively).

Plasma: It is an ionized medium which is

- ① Quasineutral
- ② exhibits a collective behaviour

that has ^{if any} - plasma is a state of a very matter state. Could be:

- ① Ionized
- ② Quasineutral.
- ③ Behave collectively.

- plasma gas of
- liquid of
- solid of

* Ionization: net charge $\neq 0$

Neutral gas
+ + + +
- - - -

$Q = 0$: (net charge $\Delta Q = 0$)

Ionized gas

+ + + + +
- - - - -

$Q \neq 0$ (net charge $\Delta Q \neq 0$)

$Q_+ > Q_-$

$Q_- > Q_+$

- It is a process by which an atom or molecule acquires a negative or positive charge by gaining or losing energy.

$E_{\text{ion}} = 13.6 \text{ eV/atom}$
 $E_0 = 12.5 \text{ eV/atom}$

* Quasineutrality: (locally)

Neutral gas

$\Delta Q = 0$

Ionized gas

$Q \neq 0$

Plasma gas

$\Delta Q \approx 0$

Charge imbalance locally

- For an ionized gas to be plasma

$E \approx 0 \Rightarrow$

$Q_+ \approx Q_-$
 $\Delta Q \approx 0$

Saha eq:

$\frac{n_i}{n_n} \sim 2.4 \times 10^{-21}$

$\frac{T^{3/2}}{n_i} e^{-U_i/k_B T}$

$\rightarrow$ for ordinary air: $\frac{n_i}{n_n} \approx 10^{-12}$; U_i : ionization energy

③

* collective behaviour :

(i) Individual behaviour:

neutral gas

- short-range interaction.
- contact force.
- mutual collision. (one to one)
- straight line trajectory
- one time interaction



"Hard sphere collision"
short-range interaction.

- The particle interact only with one at a time.

- The particle doesn't feel the other.

- the trajectory of the particle is a straight line.

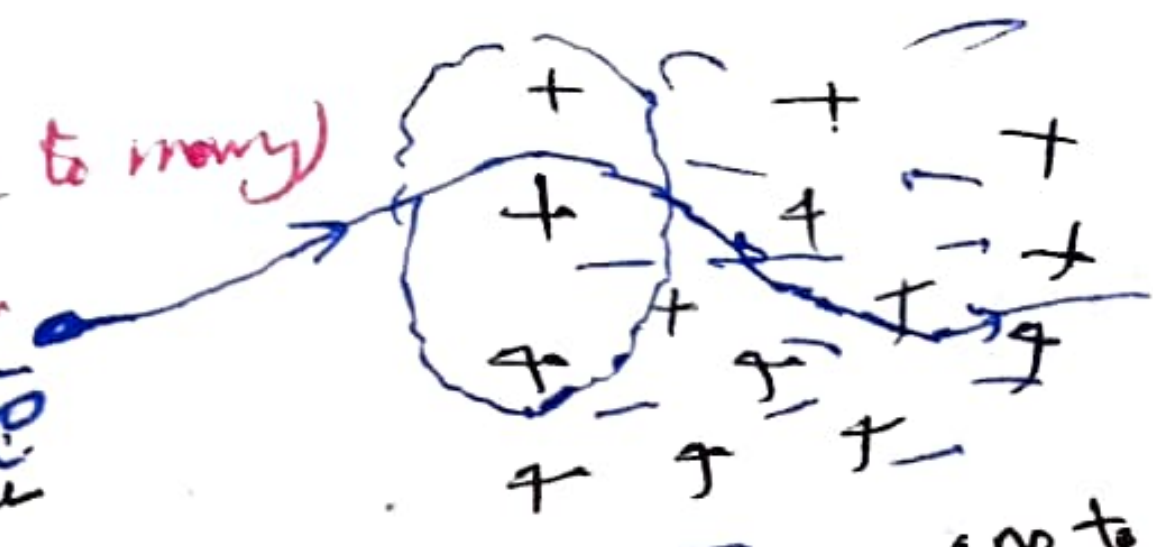
- the motion is controlled by short-range force such as collision force, contact force.

(ii) Collective behaviour:

Plasma

"Coulomb collision"
long-range interaction

- long-range interaction
- field force
- Coulomb collision (one to many)
- curvilinear trajectory.
- simultaneously interaction



inter-particle interaction
short collision
long electromagnetic
one to many

- The particle interact simultaneously with many other charged particles (not just mutual collision).

- the net force on the charge due to many other charges in certain range.

- the trajectory of the particle is curvy.

- the motion is controlled by long-range force or field force such EM force.

Neutral state

- net charge = 0
- Individual behaviour
- short range interaction: hard sphere
- contact force
- mutual collision or binary collision
- straight line trajectory

Plasma

- net charge ≈ 0
- collective behaviour
- long-range interaction: Lorentz force
- Field force
- Coulomb collision
- curvy trajectory

Plasma frequency: (natural frequency)

Plasma parameters $\rightarrow$ time scale - length λ

Assume a plasma with

(natural Inertial

$$n_e = n_0 + n_e(r, t) \quad (1)$$

$$\frac{\partial n_e(r,t)}{\partial t} + n_0 \nabla \cdot u_e(r,t) = 0 \quad (2)$$

$$\frac{\partial u_e}{\partial t} = -\frac{e}{m_e} E(r,t) \quad (3)$$

$$f = -e[n_0 + n_e] + e n_0 = -e n_e \text{ (m)} \quad (4)$$

$$\nabla \cdot E = \frac{\rho}{\epsilon_0} = -\frac{e}{\epsilon_0} n_e(r, t) \quad (5)$$

$$\nabla \cdot \textcircled{3} \Rightarrow \frac{\partial}{\partial t} \nabla \cdot \mathbf{u} = - \frac{e}{m_e} \nabla \cdot \mathbf{E}$$

$$\textcircled{2} \Rightarrow \nabla \cdot \mathbf{u}_e = \frac{1}{n_0} \frac{\partial n}{\partial t}$$

$$\frac{1}{\hbar\omega} \frac{\partial}{\partial t} \frac{\partial}{\partial t} n = -\frac{e}{m_e} \nabla \cdot E \rightarrow m \frac{\partial^2 n}{\partial t^2} = -\frac{e n_0}{m_e} \nabla \cdot E$$

from (5) $\Rightarrow \frac{\partial^2 n_e}{\partial t^2} = -\frac{n_0 e}{m_e} \cdot -\frac{e}{\epsilon_0} n_e(r, t)$

- Restoring force $= -e^2 n_0 / \epsilon_0 x = -kx$

- Inertia: m_e $\frac{\partial^2 m_e}{\partial t^2} = - \left(\frac{n_0 e^2}{\epsilon_0 m_e} \right) m_e(r, t)$

$$\frac{\partial^2 n_e}{\partial t^2} + \omega_p^2 n_e(r, t) = 0$$

where $\omega_p^2 = \frac{n_0 e^2}{\epsilon_0 m}$ $\Rightarrow \omega_p = \sqrt{\frac{n_0 e^2}{\epsilon_0 m}}$

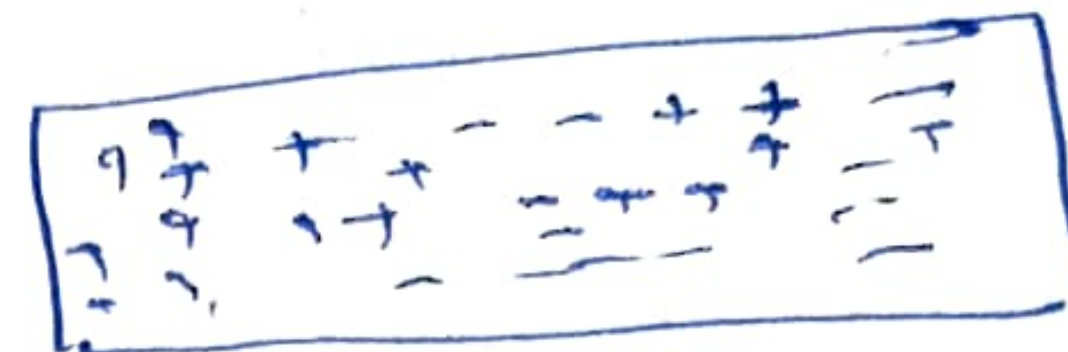
$$\omega_p = \left(\frac{e^2 n_0}{\epsilon_0 m} \right)^{1/2} \approx C \sqrt{n_0} \quad \frac{\text{rad}}{\text{sec}}$$

$$n_e(t) = n_e(r) e^{-i\omega_p t}$$

$$n_e(t) = n_e(r) \cos(\omega_p t)$$

Assume a plasma consists of electrons with charge $(-e)$ and mass m_e and same number of ions (e) and mass m_i : $n_i \approx n_e \approx n_0$
 $T_i \approx T_e \approx T$ - E

$$T_{\text{in}} T_{\text{e}} \approx T \cdot \frac{E}{\hbar}$$



- Restoring force: $E \propto \Delta x$

- Inertia: man

$\epsilon'' n_0 / s_0$ ← $F \equiv K \nabla \cdot E$

$$K \propto \frac{1}{\Sigma}$$

$+ \ominus + \ominus + \ominus + \ominus$
 $+ \ominus + \ominus + \ominus + \ominus$
 $+ \ominus + \ominus + \ominus + \ominus$

 570 K P₂O₅

$$F = - \nabla \cdot \mathbf{E}$$

Handwritten practice of the letter 't' in various orientations and sizes.

$\rightarrow$ static frequency - eno V.E
 $\rightarrow \sim \frac{ne^2}{\sigma_m}$
 $\hookrightarrow$ electrical conductivity

$$\omega_p = \sqrt{\frac{k}{m}}$$

$$K = \frac{en^2}{\epsilon}$$

medium count

$$58 \sqrt{\text{ne}(\text{cm}^{-3})} \frac{\text{rad}}{\text{fc}}$$

$$f_{\text{pe}} \approx 9 \sqrt{h e (n_0^3)} \text{ Hz}$$

Curve 10/10" width

wpv 1.8 x 10⁻¹³
f ~ 285 t13

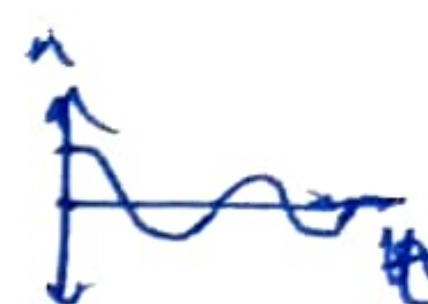
$\tau_p = \frac{1}{\omega_p}$: plasma time response to external perturbation. microwave.
~ 5-6 ps

$\omega_p \propto n_0$
 $\propto e^2$
 $\propto \frac{1}{\epsilon_0}$
 $\propto \frac{1}{m}$

$\tau_p \propto \frac{1}{n_0} \rightarrow$ more particles less time to response to E .
 $\propto \frac{1}{e^2} \rightarrow$ a chase of $\sim \sim \sim$ to E .
 $\propto \epsilon_0 \rightarrow$ high permittivity more time to E .
 $\propto m \rightarrow$ heavy mass $\sim \sim \sim$

* Collision Frequency (ν):

$\frac{d^2 n}{dt^2} + \nu \frac{dn}{dt} + \omega_p^2 n = 0$; $\nu = \frac{b}{n_{om}} = \frac{b}{f_m} \cdot \frac{\text{collision factor}}{\nu t}$ $\nu = \frac{b}{m}$

(i) light damping ($\omega_p \gg \nu$) $\Rightarrow n(r,t) = n(r) e^{-\nu t} \cos(\omega_p t)$ 

(ii) critical ($\omega_p = \nu$)  $e^{-\frac{\nu}{2} t}$

(iii) heavy: $\nu \gg \omega_p$: $n_{crit} = n(r) e^{-\frac{\nu}{2} t}$ 

$\rightarrow$ So, we in plasma interested in light damping ($\nu \ll \omega_p$).

* Thermal pressure:

Collision $\rightarrow$ with neutral
 $\rightarrow$ with charges: Coulomb collision
 $\hookrightarrow$ deflection of charges with 90° when it pass near the charges.

$\frac{\partial^2 n_e}{\partial t^2} + \nu \frac{\partial n_e}{\partial t} + \nu_{the}^2 \nabla^2 n_e + \omega_p^2 n_e = 0$

$\nu_{the} = \frac{k T_e}{m_e}$

(i) $\nu \rightarrow 0, \omega_p \rightarrow 0$

$\frac{\partial^2 n_e}{\partial t^2} + \nu_{the}^2 \nabla^2 n_e = 0$

* for ionosphere $n_e \approx 10^{12} \text{ m}^{-3}$
 $\omega_{pe} \sim 56 \sqrt{n_e} \approx 6 \times 10^7 \text{ rad/sec}$

(ii) $\nu \rightarrow 0, \nu_{the} \rightarrow 0$

$\frac{\partial^2 n_e}{\partial t^2} + \omega_p^2 n_e = 0$

$f_{pe} \approx 10 \text{ MHz} = \frac{\omega_{pe}}{2\pi}$

(iii) $\nu_{the} \rightarrow 0$

$\frac{\partial^2 n_e}{\partial t^2} + \nu \frac{\partial n_e}{\partial t} + \omega_p^2 n_e = 0$

(iv) $\frac{1}{\nu} \frac{\partial^2 n_e}{\partial t^2} + \frac{\partial n_e}{\partial t} + \frac{\nu_{the}^2}{\nu} \nabla^2 n_e + \frac{\omega_p^2}{\nu} n_e = 0$

(6) $\frac{1}{\nu} \frac{\partial^2 n_e}{\partial t^2} + \frac{\partial n_e}{\partial t} - D_e \nabla^2 n_e + \frac{\omega_p^2}{\nu} n_e = 0$; $D = \frac{k T_e}{m_e \nu} = \frac{\text{mean free path}^2}{\text{collision time}}$
 $\hookrightarrow$ Diffusion coefficient.

$V = ME - \frac{D}{n} \frac{\partial n}{\partial x}$

$D \sim \frac{(dx)^2}{dt} \sim \frac{\text{Area}}{\text{time}}$

$\mu = \frac{q}{m \nu}$

$\frac{(dx)^2}{dt}$

Debye length

1) Test charge in Vacuum:

- The electrostatic potential due to Q (Coulomb's potential)

$$\nabla \cdot E = \frac{Q}{\epsilon_0}$$

$$\nabla \cdot E = \frac{\rho}{\epsilon_0}$$

$$\nabla \cdot B = 0$$

$$\nabla \times E = -\frac{\partial B}{\partial t} \Rightarrow \nabla \times (\nabla \phi) = 0, E = -\nabla \phi$$

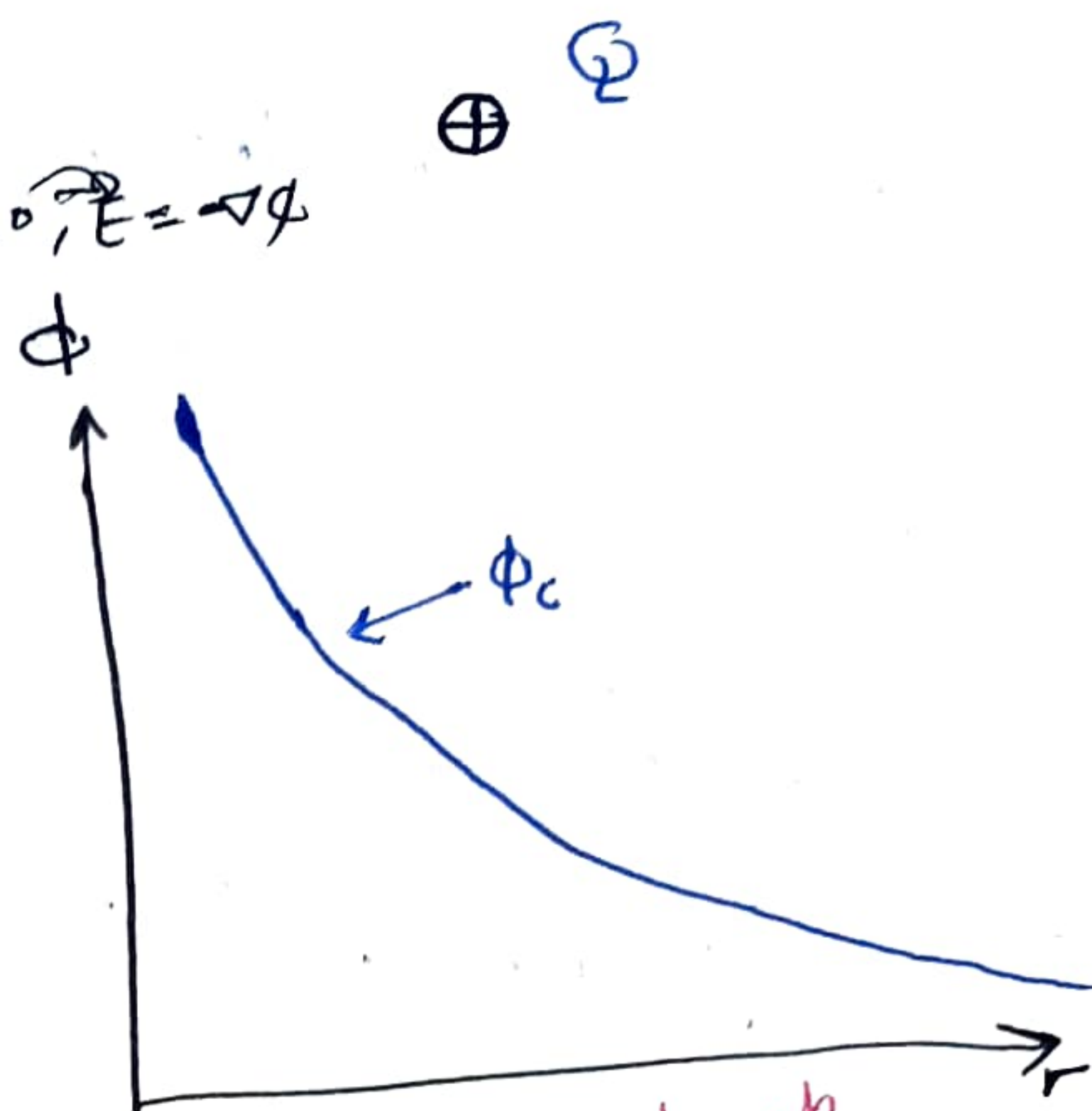
$$\nabla \times B = \mu_0 (J + \epsilon_0 \frac{\partial E}{\partial t})$$

$$E = -\nabla \phi$$

$$\nabla^2 \phi = -\frac{Q}{\epsilon_0} \Rightarrow \frac{d^2 \phi_c(r)}{dr^2} = -\frac{Q}{\epsilon_0}$$

$$\phi_c(r) = \frac{1}{\epsilon_0} \frac{1}{4\pi r} Q$$

$\phi_c \sim \frac{1}{r}$
Circumference of sphere.



Limit $\phi_c(r) \rightarrow 0$

$r \rightarrow \infty$ radius of effect (influence)

(ii) Test charge in plasma:

$$\nabla^2 \phi = -\frac{Q}{\epsilon_0} + \frac{\rho}{\epsilon_0}; \rho = en_0 e^{\frac{e\phi}{k_B T}} \approx en_0 (1 + \frac{e\phi}{k_B T} + \dots)$$

$$\rho \approx \frac{e^2 n_0}{\epsilon_0 k_B T} \phi$$

spatial simple harmonic motion
oscillation in space.

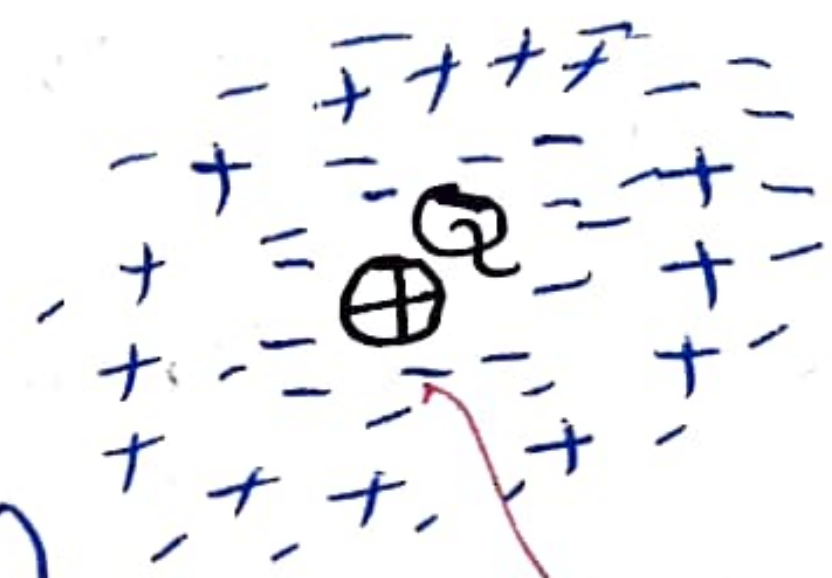
$$\nabla^2 \phi - (\frac{e^2 n_0}{\epsilon_0 k_B T}) \phi = -\frac{Q}{\epsilon_0}$$

$$\nabla^2 \phi - k_D^2 \phi = -\frac{Q}{\epsilon_0}; k_D^2 = \frac{e^2 n_0}{\epsilon_0 k_B T}$$

$$\phi(r) = \phi_c e^{-\frac{r}{\lambda_D}}$$

$$\phi(r) = \phi_c e^{-\frac{r}{\lambda_D}}$$

at $r \rightarrow \infty$ $\phi = 0$
at $r = 0$ $\phi = \phi_c$



polarization effect

for negative charges

Restoring force

spatial frequency

$$k_D^2 = \frac{e^2 n_0}{\epsilon_0 k_B T}$$

create displacement

$$\lambda_D = \frac{\epsilon_0 k_B T}{e^2 n_0}$$

$$\phi_D \sim \frac{1}{r} e^{-r}$$

- At $r \gg \lambda_D$

$$\phi \rightarrow 0$$

* physical mechanisms

- When we put the positive test charge in the plasma, the positive charge polarizes (attracts) the negative charge, i.e. negative charges polarize toward the positive charge and shield its effect. So, a charged plasma interacts effectively only with particles situated at distance less than one Debye length (λ_D) and it has a negligible influence on particles lying at distances greater than one Debye length.

$$\phi(r) \propto \frac{1}{r} e^{-r/\lambda_D}$$

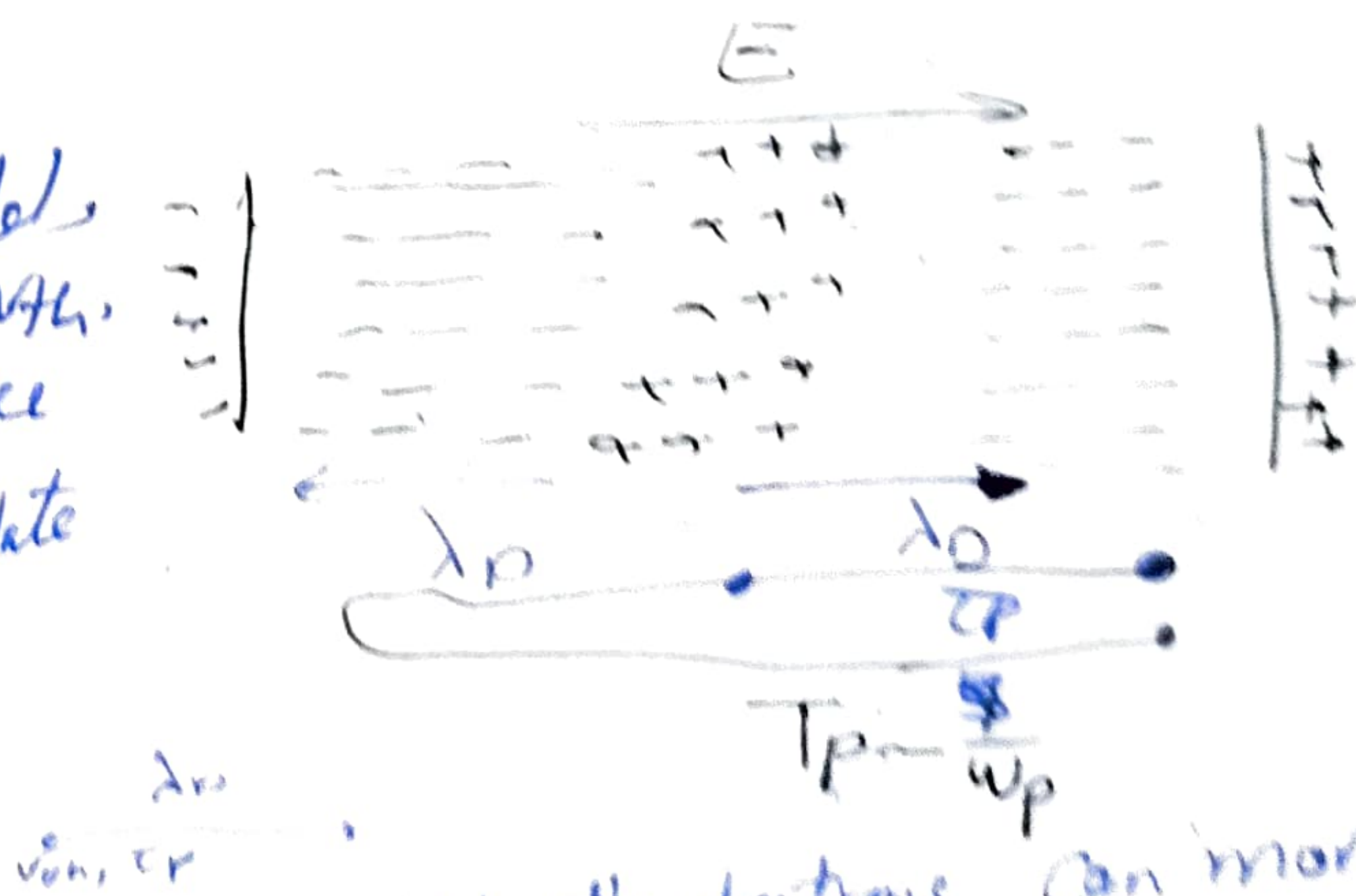
(i) $r \ll \lambda_D$: $\phi(r) \rightarrow \phi_0(r)$: charge in vacuum.

(ii) $r \gg \lambda_D$: $\phi(r) \rightarrow 0$: charge in plasma.

(iii) $r = \lambda_D$: $\phi(r) \sim \frac{\phi_0}{e} \sim 0$.

* Applying an Electric Field:

- When we apply an electric field with speed v_{th} , the plasma will displace a distance λ_D ~~in the~~ then it will oscillate with frequency ω_p .



(1) λ_D is the maximum distance over which the electrons can move with respect to the ions, against the electrostatic force with its thermal energy.

(2) the distance at that electrons move it at λ_D with velocity v_{th} (maximum displacement of oscillation).

(3) the distance beyond it, the ions have no influence on other charges.

(2)

(3)



$$D = \sqrt{\frac{e^2 n_0}{\epsilon_0 k_B T}}$$

$$\lambda_D = \sqrt{\frac{\epsilon_0 k_B T}{e^2 n_0}} \approx 6.9 \sqrt{\frac{T_e (K)}{n_0 (cm^{-3})}} \text{ cm}$$

$$\approx 7340 \text{ m} \sqrt{\frac{T_e (eV)}{n_e (m^{-3})}} \propto \sqrt{T_e} \propto \frac{1}{\sqrt{n_e}}$$

* Typical laboratory plasma:

$$T_e \sim 1 \text{ keV} = 10^3 \text{ eV}$$

$$n_e \sim 10^{19} \text{ m}^{-3} ; n = 10^{13} \text{ cm}^{-3} \text{ in vacuum.}$$

$$\lambda_D \approx 10^{-4} \text{ m} \approx 10^{-2} \text{ cm} ; \omega_{pe} \sim 1.8 \times 10^{11} \text{ rad/sec}, f_{pe} \sim 28 \text{ GHz (microwave).}$$

$$Dx \approx n^{-1/2} \approx 0.5 \times 10^{-6} \text{ m} \approx 0.5 \text{ } \mu\text{m}$$

$$a_0 \text{ (Bohr radius)} \approx 10^{-10} \text{ m}$$

$$U_{the} \sim 3 \times 10^7 \text{ m/sec.}$$

$$10^{-8} < 10^{-5} < 10^{-2} < 1$$

$$a_0 < Dx < \lambda < L$$

* Physical interpretation:

$$U_{the} = \omega_p \lambda_D ; \text{ thermal velocity.}$$

- Debye length:

① The length at which the potential $\phi = 0$

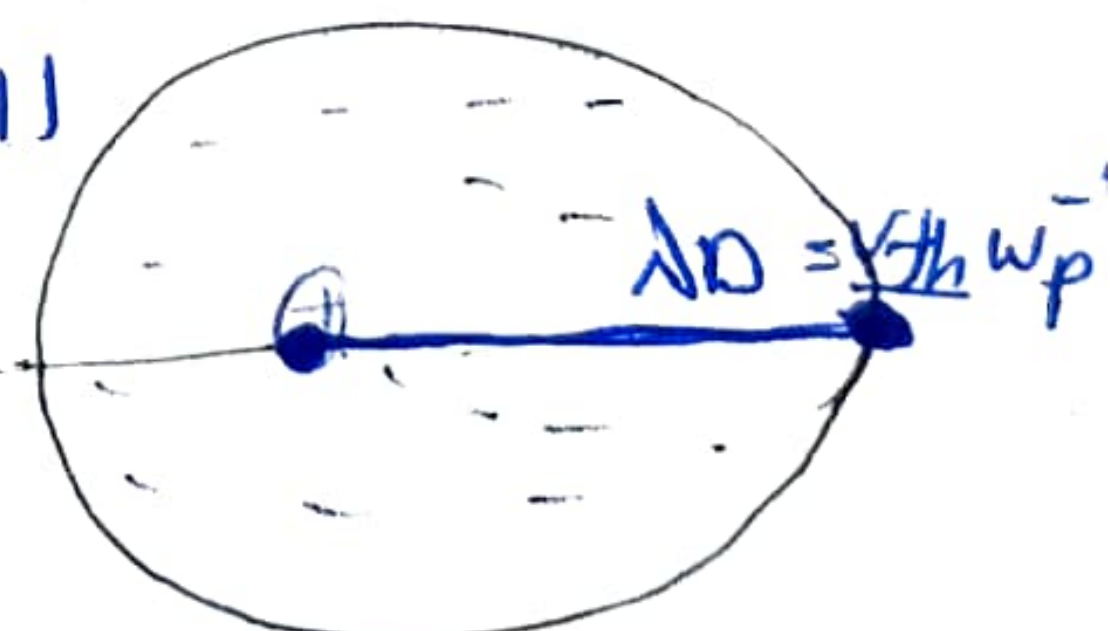
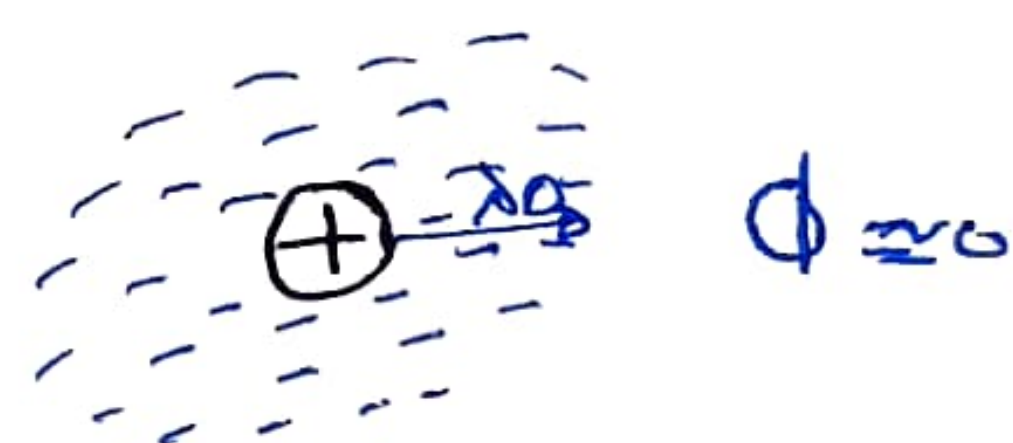
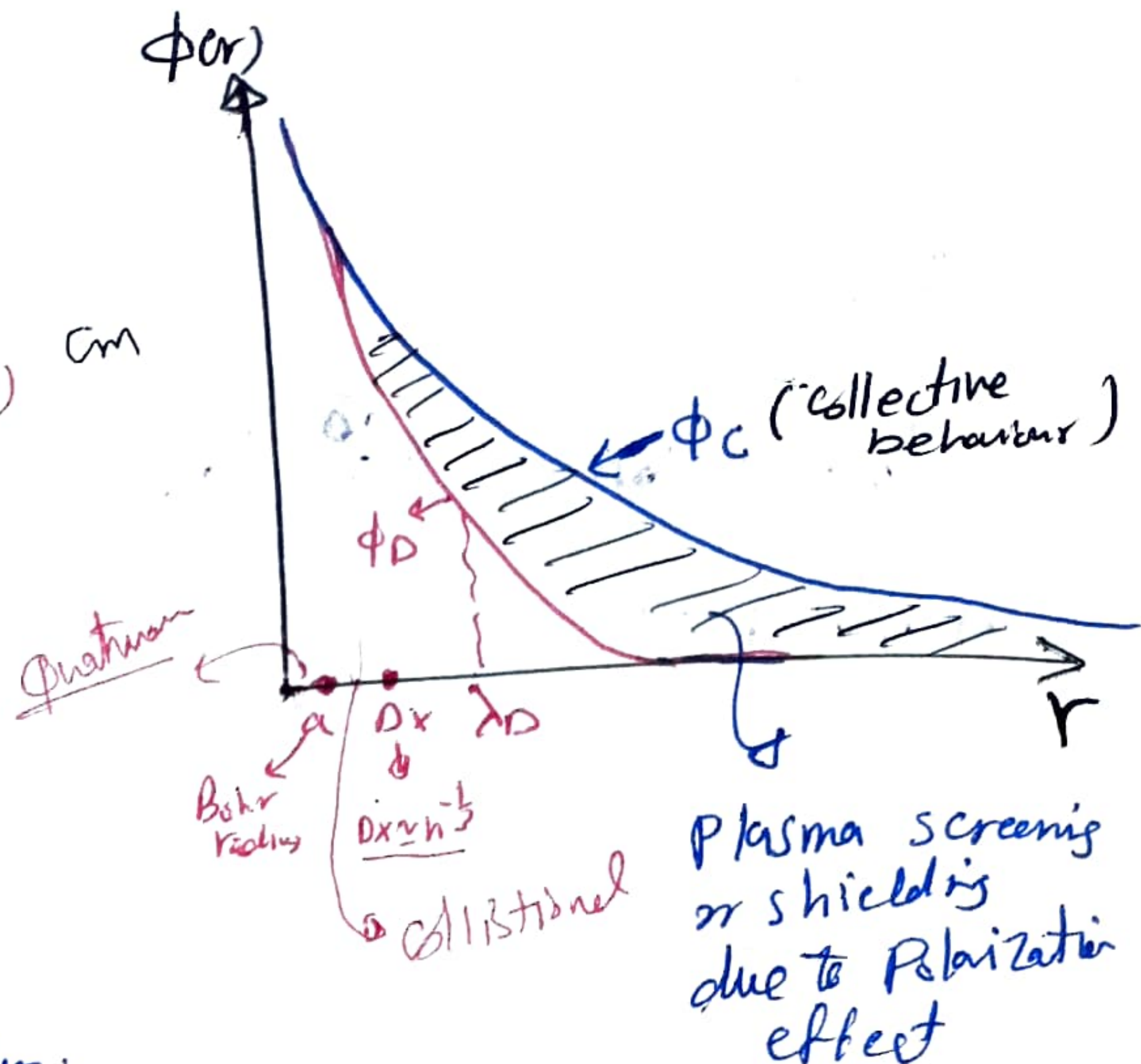
i.e.: if there is a charge at $r > \lambda_D$, it will not affect by ϕ .

② The length of shielding a test charge.

③ $\sim \sim$ the test charge interacts collectively with other charges.

④ The distance that the charge moves at $t \approx \frac{1}{\omega_p}$ with velocity V_{th} as V_{th} increases, V_{th} increases, so the charge will move further distance λ_D until it stops.

⑤ the distance over which, the plasma is neutral.



* How many particles that feels the collective behaviour (Coulomb potes) $\neq$ the particle inside the Debye sphere.

$$N_D = n \lambda_D^3 \approx 10^9 (10^{-9})^3 \approx 10^0 \neq \frac{\#}{m^3}$$

$n \lambda_D^3$: plasma parameter.

* Note:

$$\frac{\partial^2 n e}{\partial x^2} + \omega_p^2 m = 0 \quad ; \quad \omega_p^2 = \frac{e^2 n_0}{\epsilon_0 m e} \quad ; \quad n \sim e^{-i\omega t} \quad ; \quad + \equiv -i$$

$$\frac{\partial^2 \phi}{\partial x^2} - k_D^2 \phi = 0 \quad ; \quad k_D^2 = \frac{e^2 n_0}{\epsilon_0 k_B T_e} \quad ; \quad \phi \sim e^{-k_D x} \quad ; \quad - \equiv -1 \quad ; \quad \text{Response}$$

→ temporal frequency:

$$\omega_p^2 = \frac{\text{Restoring force to neutral}}{\text{Perturbation (Inertia)}} = \frac{e^2 n_0 / \epsilon_0}{m} \quad ; \quad \omega_p \sim \frac{1}{T_p} \quad ; \quad T_p: \text{Response time}$$

* Spatial frequency:

$$k_D^2 = \frac{\text{Restoring force to}}{\text{Inertia}} = \frac{e^2 n_0 / \epsilon_0}{k_B T_e} \quad ; \quad k_D \sim \frac{1}{\lambda_D} \quad ; \quad \lambda_D: \text{Response length}$$

* Restoring force: to restore the neutrality or electrical equilibrium

* Inertia: try to keep the non-equilibrium situation.

→ charge imbalances may exist only over a short distance (Debye length)
or for a short period of time (inverse of plasma frequency).

plasma criteria : (Quantitatively)

The conditions for any ionized medium to be considered a plasma

① Quasineutrality:

- From Gauss's law:

ϵ

$$\frac{1}{\epsilon_0} \nabla \cdot \mathbf{E} = \nabla^2 \phi = n_e - n_i$$

$$\frac{\lambda_D^2}{L^2} \phi = n_e - n_i$$

for $\lambda_D \ll L : LHS = 0$

$$\therefore n_e - n_i \approx 0 \rightarrow n_e \approx n_i$$

So, the condition for quasineutrality

$$L \gg \lambda_D$$

plasma size: L

high temperature and low dense plasma

- Typical laboratory plasma

$$L \approx 10 \text{ cm} \approx 0.1 \text{ m} \approx 10^3 \lambda_D$$

* Conditions:

① collective behavior:

$$\Gamma \gg 1; N_D \gg 1$$

necessary condition

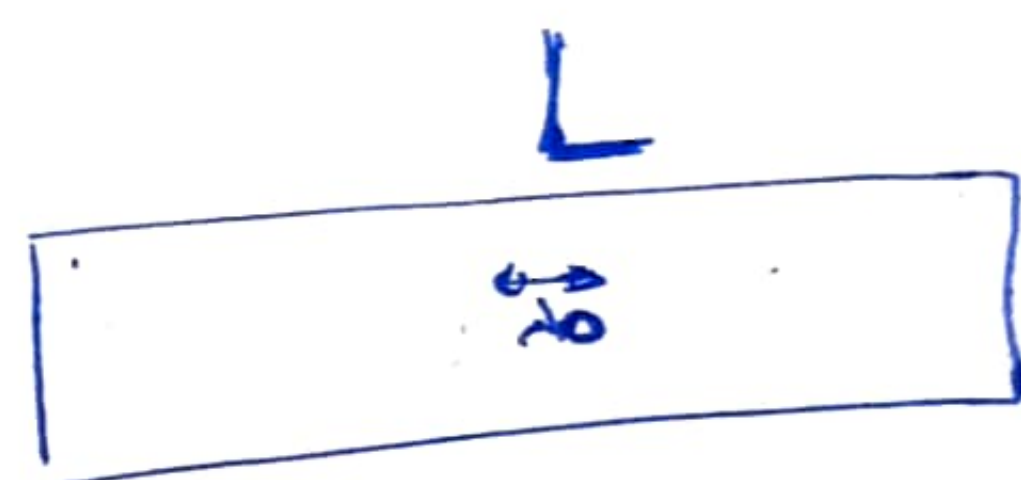
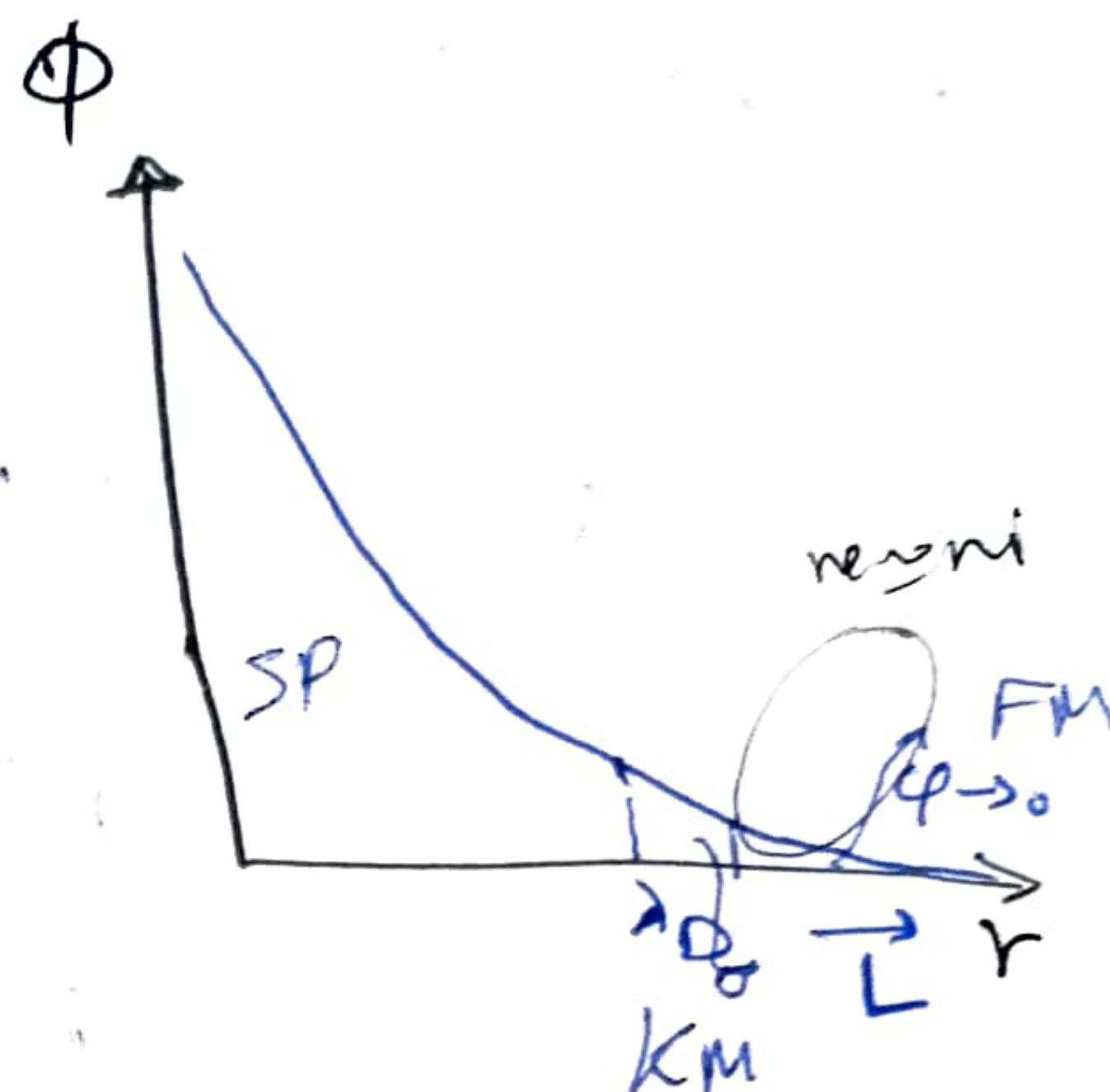
② quasineutrality:

sufficient ~

③ $\omega_p \gg \omega_c$

$$\omega_p \gg \omega_c$$

sufficient



② Collective behaviour:

- ~~that~~ the mean interparticle distance, collective interaction ~~of~~ of charged particles dominate over binary interaction.

$$\frac{e\phi}{k_B T_e} \approx \frac{1}{4\pi(n\lambda_D^3)^{2/3}} \ll 1$$

$$n\lambda_D^3 \gg 1$$

$$\boxed{N_D \gg 1}$$

$$\Gamma = \frac{e\phi}{k_B T_e} \ll 1 \rightarrow N_D \gg 1$$

$$\Gamma = \frac{e\phi}{k_B T_e} \ll 1 \text{ (weakly coupled)} \rightarrow \text{one to one interaction}$$

- It means you need so many particles inside the Debye sphere to interact collectively.

- Typical laboratory plasma:

$$N_D \sim 10^6 \gg 1$$

- there is a million particles inside the Debye sphere.

$$\therefore L \approx 10^3 \lambda_D \Rightarrow (N_D)^{2/3} \approx 10^{18} \text{ m}^{-3}$$

- there is $10^{19} \frac{1}{\text{m}^3}$ in L .

→ Another meaning (Coupling parameter) Γ :

$$\Gamma = \frac{E_c}{E_{th}} \ll 1$$

↳ plasma parameter

$$\frac{\text{Thermal fluctuation}}{\text{Thermal energy density}} = \frac{\epsilon_0 |E|^2}{n e T_e} \sim \frac{1}{n \lambda_D^3} \ll 1$$

$$\boxed{n \lambda_D^3 \gg 1}$$

So the thermal fluctuations level is so small compared to the thermal energy density

Weakly collision (weakly coupled):

- The long-range interaction force is dominate over the short-range interaction force or binary collision force.
- the plasma oscillate many times before, the collision happen.

- the time response of the plasma to the long-range force is smaller than the time response to the collision:

$$\Lambda = \frac{4}{3} \frac{v_{th}^3}{\omega_p^3} n$$

$$\nu \sim \frac{\ln \Lambda}{\Lambda} \omega_p$$

$$\nu_c \sim \frac{\omega_p}{(n \lambda_D)^3}$$

diffuse
low dense plasma
high temperature

(i) $\omega_p \gg \nu_c$: weakly coupled plasma: $\Lambda \gg 1$
collision do not interfere with plasma oscillation

dense
low temperature

(ii) $\nu_c \gg \omega_p$: strongly coupled plasma: $\Lambda \ll 1$
collisions effectively prevent plasma oscillation

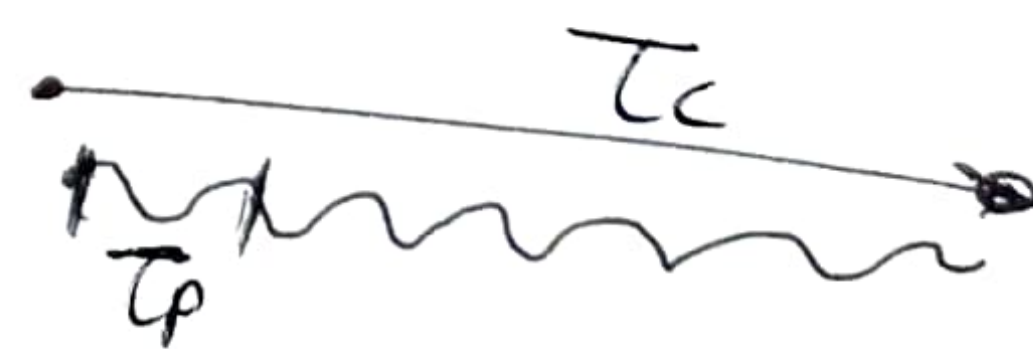
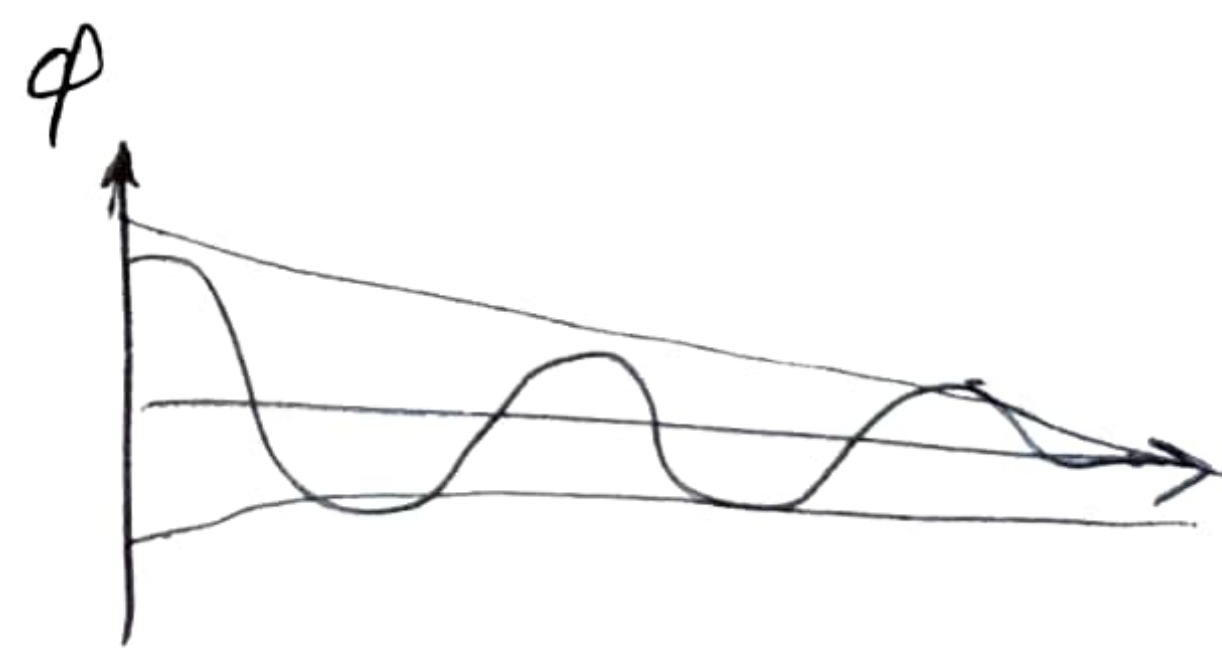
$$\therefore (n \lambda_D)^3 \gg 1$$

(i) $\nu_c \ll \omega_p$

(light-damping)

(ii) $\tau_c \gg \tau_p$

(iii) $\lambda_{mfp} \gg L$: collisions plasma.



$$\lambda_{mfp} = \frac{v_{th}}{\nu_c} \approx v_{th} \tau_c$$

ν_c : is the inverse of the typical time needed for enough collision to occur that the particle trajectory is deviated through 90° or 90° scattering rate.



* Plasma in a tokamak:

$$n_e \approx 10^{14} \text{ cm}^{-3}, T_e = 1 \text{ keV}, L = 20 \text{ cm}$$

So, the plasma parameters

$$v_{the} \approx 2 \times 10^9 \sqrt{\frac{2T_e}{m_e c}} \approx 2 \times 10^9 \text{ cm/sec}$$

$$\omega_{pe} \approx 5.6 \times 10^{11} \text{ sec}^{-1}$$

$$\lambda_{De} \approx 0.003 \text{ cm}$$

$$L = 20 \text{ cm}$$

$$ND = n_e d_{De}^3 \approx 3 \times 10^6$$

$$a_0 (\text{Bohr radius}) \approx 10^{-8} \text{ cm}$$

$$Dx \approx n^{-1/3} \approx 10^{-5} \text{ cm} \quad \text{--- Interparticle distance.}$$

⇒ plasma conditions

① No quantum effect

$$a_0 \ll Dx \ll \lambda_D \ll L$$

$$\text{cm } 10^{-8} \ll 10^{-5} \text{ cm} \ll 3 \times 10^{-3} \text{ cm} \ll 20 \text{ cm}$$

$$② ND \gg 1$$

$$① \lambda_D < Dx : \text{classical}$$

$$② \lambda_D < a_0 : \text{quantum}$$

Category Plasma classifications (scale):

- ① Laboratory plasma: (earth-scale length scale).
- ② Space plasma: (Solar-system length scale).
- ③ Astrophysics plasma: (Galaxy-length scale).
- ④ Cosmology plasma: (universe-length scale).

* Plasma category: Classification by n_e, T_e $n^{-1/3} \ll \lambda_D \ll L$

① classical plasma: $L \gg \lambda_{DB}$ OR $T_{th} \gg T_F$.
 $\lambda_D \ll d$; in gold: $n \approx 10^{23}$

② Quantum plasma: $L \sim \lambda_{DB}$ OR $T_F \gg T_{th}$; $E_T < 2.75 E_F$
 $\lambda_D < a_0$

$\log T (eV)$

10
8
7
6
5
4
3
2
1

nebulas
 Solar wind
 Interstellar

Russias
 Solar corona
 Neon lamps

$\log(T_e)$
 15
10
5

Solar core
 Solar corona
 Ionosphere
 Fusion
 Gas discharge

$N_D = 10^{10}$
 $N_D = 1$

Solar liquid gas

$\log(n)$

15

"Parameter Space of Plasma"

③ Relativistic plasma:

$$kT > m_e c^2$$

④ Collisional plasma:

$$D \gg \lambda_D \gg \lambda_{on}^2, \lambda_{on}^2 < 1$$

interparticle dist

* Plasma applications:

* Fusion:

Lawson criterion

$$n\tau > 1.7 \times 10^{14} \frac{\text{sec}}{\text{cc}}$$

low density

cc: cm³

long confinement time

MCF: $n = 10^{14} / \text{cc}$

high density

$$\rho\tau = 1 \text{ sec}$$

start - confinement time

ICF: $n = 10^{25} / \text{cc}$

$$\rho\tau = 10^{11} \text{ sec}$$